

Homogeneous quaternionic Kähler structures on eight-dimensional non-compact quaternion-Kähler symmetric spaces^{*}

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Abstract

For each non-compact quaternion-Kähler symmetric space of dimension 8, all of its descriptions as a homogeneous Riemannian space (obtained through the Witte's refined Langlands decomposition) and the associated homogeneous quaternionic Kähler structures are studied.

1 Introduction and Preliminaries

Introduction. As it is well-known, quaternion-Kähler symmetric spaces were classified by Wolf [21, Th. 5.4]. The non-compact ones are moreover a particular case of Alekseevsky spaces ([1, p. 338], [8, Th. 2.28]). On the other hand, a theoretical-representation classification of homogeneous quaternionic Kähler structures was given by Fino [9, L. 5.1]. Then, a classification by real tensors was obtained in [6, Th. 3.15], where furthermore a characterization of the quaternionic hyperbolic space in terms of such structures was given.

The aim of this paper is to find all the homogeneous Riemannian structures associated to every homogeneous description (obtained through Witte's refined Langlands decomposition [20, Th. 1.2]) of each 8-dimensional non-compact quaternion-Kähler symmetric space, and to determine their types as homogeneous quaternionic Kähler structures.

Non-compact quaternion-Kähler spaces are found in the formulation of the coupling of matter fields in $N = 2$ supergravity as target spaces of some non-linear sigma models (Bagger-Witten [4, p. 7], Cecotti [7, p. 32], de Wit-Van Proeyen [19, p. 26]; see also Fré *et al.* [10, p. 65] and references therein). Since Fino's classification is natural, it seems reasonable the existence of links between the classification of homogeneous quaternionic Kähler structures and a possible

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classification of some physical structures and models. This would, in a sense, be similar to a recent result by Meessen [16, Th. 2]: A connected homogeneous Lorentzian space admits a non-vanishing degenerate homogeneous Lorentzian structure in Tricerri-Vanhecke's class $\mathcal{T}_1 \oplus \mathcal{T}_3$ (see [18, p. 41], [11, Th. 2.1], [17, p. 562]) if and only if it is a singular homogeneous plane wave. Moreover, we give in the present paper some candidates to such target manifolds.

As for the contents, we first establish (Theorem 2.3) a result related to Witte's Theorem on closed co-compact subgroups of a semisimple Lie group with finite center ([20, Th. 1.2]), which gives necessary and sufficient conditions for a connected closed co-compact subgroup of such a group to act transitively (see also Gordon-Wilson [12, Th. 6.9]). We then furnish explicitly (Proposition 2.4), for the present case of symmetric spaces, the coefficients appearing in the symmetries of homogeneous structures characteristic of the quaternionic Kähler case (that is, equations (1.3)). Further, for each one of spaces under study, i.e. the complex hyperbolic Grassmannian $SU(2, 2)/S(U(2) \times U(2))$, the quaternionic hyperbolic space $\mathbb{H}H(2) = Sp(2, 1)/(Sp(2) \times Sp(1))$, and the exceptional space $G_{2(2)}/SO(4)$, we first give the corresponding quaternionic structure. This is associated to a natural structure of quaternion-Hermitian vector space on the Lie algebra $\mathfrak{a} + \mathfrak{n}$ of the solvable factor AN of an Iwasawa decomposition of its respective full connected group of isometries. According to Alekseevsky [1, Th. 1.1] (see also Heber [13, Cor. 5.5]), its metric is unique up to isometry and scaling.

Then, we obtain all the descriptions of each space and the corresponding homogeneous quaternionic Kähler structures. These structures are associated to the reductive decompositions defined via the Witte's refined Langlands decomposition of the identity component of each parabolic subgroup of the full connected isometry group. It turns out that each parabolic subgroup provides a homogeneous description which, although is the simplest possible one, suffices to obtain all such homogeneous structures. We exhibit the corresponding structures in Propositions 3.1, 3.2, 3.3, and their types in Theorem 3.4.

We have studied the twelve-dimensional case in [5].

Homogeneous quaternionic Kähler structures. Let (M, g) be a connected, simply-connected, and complete Riemannian manifold. Ambrose and Singer [3, p. 656] gave a characterisation for (M, g) to be homogeneous in terms of a $(1, 2)$ tensor field S , usually called (Tricerri-Vanhecke [18, Def. 2.1]) a homogeneous Riemannian structure: Let ∇ be the Levi-Civita connection of g and R its curvature tensor. Then the manifold is homogeneous Riemannian if and only if there exists a $(1, 2)$ tensor field S such that the Ambrose-Singer equations

$$(1.1) \quad \tilde{\nabla}g = 0, \quad \tilde{\nabla}R = 0, \quad \tilde{\nabla}S = 0,$$

where $\tilde{\nabla} = \nabla - S$, are satisfied.

Let (M, g, v) be a quaternion-Kähler manifold, v denoting the distinguished rank-three subbundle of the bundle of $(1, 1)$ tensor fields on M . Then (M, g, v) is said to be a *homogeneous quaternion-Kähler space* if it admits a transitive

group of isometries (Alekseevsky-Cortés [2, Th. 1.1]). The canonical 4-form Ω is known to be globally defined and we have

Corollary 1.1. (To KIRIČENKO [14, Th. 1]) *A connected, simply-connected, and complete quaternion-Kähler manifold (M, g, ν) is homogeneous if and only if there exists a tensor field S of type (1, 2) on M satisfying*

$$(1.2) \quad \tilde{\nabla}g = 0, \quad \tilde{\nabla}R = 0, \quad \tilde{\nabla}S = 0, \quad \tilde{\nabla}\Omega = 0,$$

where $\tilde{\nabla} = \nabla - S$.

Such a tensor S is called a *homogeneous quaternionic Kähler structure* on M . The condition $\tilde{\nabla}\Omega = 0$ can be replaced by the set of equations

$$(1.3) \quad \begin{aligned} S_X J_1 &= \theta^3(X)J_2 - \theta^2(X)J_3, \\ S_X J_2 &= \theta^1(X)J_3 - \theta^3(X)J_1, \\ S_X J_3 &= \theta^2(X)J_1 - \theta^1(X)J_2, \end{aligned}$$

for certain differential 1-forms θ^a , $a = 1, 2, 3$, where $\{J_1, J_2, J_3\}$ is a local basis of ν satisfying the conditions $J_a^2 = -I$, $J_a J_b = -J_b J_a = J_c$, for each cyclic permutation (a, b, c) of $(1, 2, 3)$. Denoting $S_{XYZ} = g(S_X Y, Z)$, one has

$$(1.4) \quad S_{XJ_1 Y J_1 Z} - S_{XYZ} = \theta^3(X)g(J_2 Y, J_1 Z) - \theta^2(X)g(J_3 Y, J_1 Z),$$

and the two other similar equations.

Note that the conditions (1.4) (equivalently, equations (1.3)), besides the general condition $S_{XYZ} = -S_{XZY}$, are the symmetries satisfied by a homogeneous quaternionic Kähler structure S .

A quaternion-Hermitian real vector space $(V, \langle \cdot, \cdot \rangle, J_1, J_2, J_3)$ is the model for the tangent space at any point of a quaternion-Kähler manifold. Consider the space of tensors $\mathcal{T}(V) = \{S \in \otimes^3 V^* : S_{XYZ} = -S_{XZY}\}$ and its vector subspace \mathcal{V} of tensors satisfying (1.4) with $\langle \cdot, \cdot \rangle$, $\theta^a \in V^*$. Then $\mathcal{V} = \check{\mathcal{V}} + \hat{\mathcal{V}}$, where

$$\begin{aligned} \check{\mathcal{V}} &= \{ \Theta \in \otimes^3 V^* : \Theta_{XYZ} = \sum_{a=1}^3 \theta^a(X) \langle J_a Y, Z \rangle, \theta^a \in V^* \}, \\ \hat{\mathcal{V}} &= \{ T \in \otimes^3 V^* : T_{XYZ} = -T_{XZY}, T_{XJ_a Y J_a Z} = T_{XYZ}, a = 1, 2, 3 \}. \end{aligned}$$

The decomposition above of \mathcal{V} is orthogonal with respect to the scalar product (\cdot, \cdot) defined by $(S, S') = \sum_{r,s,t=1}^{4n} S_{e_r e_s e_t} S'_{e_r e_s e_t}$, where $\{e_r\}_{r=1,\dots,4n}$ is an orthonormal basis of V . Then $\check{\mathcal{V}}$ and $\hat{\mathcal{V}}$ decompose into two and three subspaces, respectively, giving an orthonormal sum of five subspaces which are invariant and irreducible under the action of $Sp(n)Sp(1)$, due to the following theorem. Let E denote the standard representation of $Sp(n)$ on \mathbb{C}^{2n} ; $S^3 E$ the 3-symmetric product of E ; K the irreducible $Sp(n)$ -module of highest weight $(2, 1, 0, \dots, 0)$; H the standard representation of $Sp(1)$ on \mathbb{C}^2 ; and $S^3 H$ the 4-dimensional symmetric product of H . Using brackets for real representations and with the usual notations we have

Theorem 1.2. (FINO [9, L. 5.4]) *The space $[EH] \otimes (\mathfrak{sp}(1) \oplus \mathfrak{sp}(n))$ of homogeneous quaternionic Kähler structures splits into invariant and irreducible subspaces under the action of $Sp(n)Sp(1)$ as $[EH] + [ES^3H] + [EH] + [S^3EH] + [KH]$.*

Let \mathcal{QK}_i be the i th summand in Fino's classification. Then we have

Theorem 1.3. ([6, Th. 3.15]) *If $n \geq 2$, then \mathcal{V} decomposes into the direct sum of the following subspaces invariant and irreducible under the action of $Sp(n)Sp(1)$:*

$$\begin{aligned}\mathcal{QK}_1 &= \{ \Theta \in \check{\mathcal{V}} : \Theta_{XYZ} = \sum_{a=1}^3 \theta(J_a X) \langle J_a Y, Z \rangle, \theta \in V^* \}, \\ \mathcal{QK}_2 &= \{ \Theta \in \check{\mathcal{V}} : \Theta_{XYZ} = \sum_{a=1}^3 \theta^a(X) \langle J_a Y, Z \rangle, \sum_{a=1}^3 \theta^a \circ J_a = 0, \theta^a \in V^* \}, \\ \mathcal{QK}_3 &= \{ T \in \hat{\mathcal{V}} : T_{XYZ} = \langle X, Y \rangle \vartheta(Z) - \langle X, Z \rangle \vartheta(Y) \\ &\quad + \sum_{a=1}^3 (\langle X, J_a Y \rangle \vartheta(J_a Z) - \langle X, J_a Z \rangle \vartheta(J_a Y)), \vartheta \in V^* \}, \\ \mathcal{QK}_4 &= \{ T \in \hat{\mathcal{V}} : T_{XYZ} = \frac{1}{2} (T_{YZX} - T_{ZXY} \\ &\quad + \sum_{a=1}^3 (T_{J_a Y J_a Z X} - T_{J_a Z J_a Y X})), c_{12}(T) = 0 \}, \\ \mathcal{QK}_5 &= \{ T \in \hat{\mathcal{V}} : T_{XYZ} = -\frac{1}{4} (T_{YZX} - T_{ZYX} + \sum_{a=1}^3 (T_{J_a Y J_a Z X} - T_{J_a Z J_a Y X})) \}.\end{aligned}$$

We will denote the sum of classes $\mathcal{QK}_i + \mathcal{QK}_j$ by \mathcal{QK}_{ij} , and so on. In particular, we have $\check{\mathcal{V}} = \mathcal{QK}_{12}$, $\hat{\mathcal{V}} = \mathcal{QK}_{345}$.

Remark 1.4. Let (M, g, v) be a homogeneous quaternion-Kähler space. For any given $p \in M$, the decomposition of \mathcal{V}_p given by Theorem 1.3 depends only on v_p and not on the chosen bases $\{(J_1)_p, (J_2)_p, (J_3)_p\}$, so the irreducible summands $(\mathcal{QK}_i)_p$ give well-defined bundles \mathcal{QK}_i over M . Moreover, if S is a homogeneous quaternionic Kähler structure and S_p belongs to a given invariant subspace of \mathcal{V}_p , then S_q belongs to the invariant subspace of \mathcal{V}_q of the same type at any other $q \in M$, and S is a section of the corresponding vector bundle.

2 Non-compact homogeneous spaces

Non-discrete co-compact subgroups acting transitively. Gordon and Wilson gave in [12, Th. 6.9] a theorem of characterization of the isometry groups acting transitively on non-compact Riemannian symmetric spaces. We will prove Theorem 2.3, related to Witte's Theorem 2.2, which suffices for our purposes.

Let (M, g) be a connected non-compact Riemannian manifold and G a Lie group of isometries acting transitively and effectively on M . If K is the isotropy group at any fixed point $o \in M$, then $M = G/K$. We look for the connected closed subgroups \hat{G} of G acting transitively by isometries on M . Denoting $K_{\hat{G}} = \hat{G} \cap K$, one must have $M \equiv \hat{G}/K_{\hat{G}}$, and thus other possible description of (M, g) as a homogeneous Riemannian space. Note that since G acts effectively then K is compact, and if $G = \hat{G}K$ then $G/\hat{G} \equiv K/K_{\hat{G}}$. Then we have

Lemma 2.1. (1) A subgroup \hat{G} of G acts transitively on M if and only if $G = \hat{G}K$. (2) If \hat{G} is a closed subgroup of G which acts transitively on M , then it is co-compact, that is, G/\hat{G} is compact.

The structure of the non-discrete co-compact subgroups of a connected semi-simple Lie group G with finite center was given in Witte [20, Th. 1.2]: Let \mathfrak{g} be the Lie algebra of G , \mathfrak{a} a maximal \mathbb{R} -diagonalizable subalgebra of \mathfrak{g} , Σ the set of roots of $(\mathfrak{g}, \mathfrak{a})$, and $\mathfrak{g} = \mathfrak{g}_0 + \sum_{f \in \Sigma} \mathfrak{g}_f$ the restricted-root space decomposition, with $\mathfrak{g}_0 = \mathfrak{a} + Z_{\mathfrak{k}}(\mathfrak{a})$, where $Z_{\mathfrak{k}}(\mathfrak{a})$ stands for the centralizer of \mathfrak{a} in \mathfrak{k} . Write Σ^+ for the set of positive roots with respect to a certain notion of positivity for \mathfrak{a}^* , and let Π be the set of simple restricted roots. For each subset Ψ of Π , let $[\Psi]$ be the set of restricted roots that are linear combinations of elements of Ψ . Then, the (connected) standard parabolic subgroup P_{Ψ}^0 is defined as the connected subgroup of G having Lie algebra $\mathfrak{p}_{\Psi} = \mathfrak{g}_0 + \sum_{f \in \Sigma^+ \cup [\Psi]} \mathfrak{g}_f = \mathfrak{l}' + \mathfrak{n}_{\Psi}$, where $\mathfrak{l}' = \mathfrak{g}_0 + \sum_{f \in [\Psi]} \mathfrak{g}_f = \mathfrak{l}'_{\Psi} + \mathfrak{e}'_{\Psi} + \mathfrak{a}_{\Psi}$, with \mathfrak{l}'_{Ψ} semisimple with non-compact summands, \mathfrak{e}'_{Ψ} compact reductive, \mathfrak{a}_{Ψ} the non-compact part of the center of \mathfrak{l}' , and with $\mathfrak{n}_{\Psi} = \sum_{f \in \Sigma^+ \setminus [\Psi]} \mathfrak{g}_f$ nilpotent. On the Lie group level we have [20, Th. 1.2] the *refined Langlands decomposition* $P_{\Psi}^0 = L'_{\Psi} E'_{\Psi} A_{\Psi} N_{\Psi}$, and one has the

Theorem 2.2. (WITTE [20, Th. 1.2]) Let L be a connected normal subgroup of L'_{Ψ} and E a connected closed subgroup of $E'_{\Psi} A_{\Psi}$. Then there is a closed co-compact subgroup \hat{G} of G contained in P_{Ψ} with identity component $\hat{G}^0 = LEN_{\Psi}$. Further, every closed co-compact subgroup of G arises in this way.

If $I_0(M)$ is the full connected isometry group of M , to find the homogeneous descriptions of the non-compact Riemannian symmetric space (M, g) , we will look among the co-compact subgroups \hat{G} of $I_0(M)$ given by Witte's Theorem 2.2 for those connected groups acting transitively. Note that if a Lie group acts transitively on M , so does its identity component. From Lemma 2.1 and Theorem 2.2, and writing $N = N_{\emptyset}$ when $\Psi = \emptyset \subset \Pi$, we have

Theorem 2.3. Let G be a connected semisimple Lie group with finite center and $G = KAN$ the Iwasawa decomposition. A connected closed co-compact subgroup $\hat{G} = LEN_{\Psi}$ of G acts transitively on $M = G/K$ if and only if: (a) The projections of the Lie algebra $\mathfrak{l} \subset \mathfrak{l}' = \mathfrak{g}_0 + \sum_{f \in [\Psi]} \mathfrak{g}_f$ of L to $\sum_{f \in \Sigma^+ \cap [\Psi]} \mathfrak{g}_f$ and to $\mathfrak{a}_{\Psi}^{\perp}$ are surjective, $\mathfrak{a}_{\Psi}^{\perp}$ being the orthogonal complement to \mathfrak{a}_{Ψ} in $\mathfrak{a} \subset \mathfrak{g}_0 = \mathfrak{a}_{\Psi} + \mathfrak{a}_{\Psi}^{\perp} + Z_{\mathfrak{k}}(\mathfrak{a})$. (b) The projection of the Lie algebra $\mathfrak{e} \subset \mathfrak{e}'_{\Psi} + \mathfrak{a}_{\Psi}$ of E to \mathfrak{a}_{Ψ} is surjective.

Proof. Transitivity is equivalent to $\hat{G}K = G$. Since \hat{G} and K are connected subgroups of G , this is equivalent to saying that the Lie algebra $\hat{\mathfrak{g}} + \mathfrak{k}$ of $\hat{G}K$ coincides with \mathfrak{g} . On the other hand, the Iwasawa decomposition of G gives us the vector space direct sum decomposition $\mathfrak{g} = \mathfrak{k} + \mathfrak{a} + \mathfrak{n}$. Then, if $\pi_{\mathfrak{k}}$, $\pi_{\mathfrak{a}}$, and $\pi_{\mathfrak{n}}$ denote the projections to the summands \mathfrak{k} , \mathfrak{a} and \mathfrak{n} , a necessary and sufficient condition for $\hat{\mathfrak{g}} + \mathfrak{k} = \mathfrak{g}$ to be fulfilled, is to have $\pi_{\mathfrak{a}}(\hat{\mathfrak{g}}) = \mathfrak{a}$ and $\pi_{\mathfrak{n}}(\hat{\mathfrak{g}}) = \mathfrak{n}$. Hence we have that: (1) Since $\mathfrak{n} = \sum_{f \in \Sigma^+} \mathfrak{g}_f$ and $\mathfrak{n}_{\Psi} = \sum_{f \in \Sigma^+ \setminus [\Psi]} \mathfrak{g}_f$, to get $\pi_{\mathfrak{n}}(\hat{\mathfrak{g}}) = \mathfrak{n}$, the first condition in (a) must be true, and (2) To have $\pi_{\mathfrak{a}}(\hat{\mathfrak{g}}) = \mathfrak{a}$,

as on the one hand E is a subgroup of $E'_\Psi A_\Psi$ and on the other hand $\mathfrak{a}_\Psi \subset \mathfrak{a}$, the other condition in (a) and the one in (b) must be satisfied. \square

Homogeneous Riemannian structures. Let (M, g, v) a connected simply-connected homogeneous quaternion-Kähler space. We have $M = G/K$, where $G = I_0(M)$ and K is the isotropy subgroup of G at a point $o \in M$. Let \hat{G} be a connected closed Lie subgroup of G which acts transitively on M . The isotropy group of this action at $o = K \in M$ is $H = K_{\hat{G}} = \hat{G} \cap K$. Then $M = G/K$ has also the description $M \equiv \hat{G}/H$, and $o \equiv H \in \hat{G}/H$.

Consider a reductive decomposition of the Lie algebra $\hat{\mathfrak{g}}$ of \hat{G} , that is, a vector space direct sum $\hat{\mathfrak{g}} = \mathfrak{h} + \mathfrak{m}$, where \mathfrak{h} is the Lie algebra of H and $\text{Ad}(H)\mathfrak{m} \subset \mathfrak{m}$. Since H is connected, the condition $\text{Ad}(H)\mathfrak{m} \subset \mathfrak{m}$ is equivalent to $[\mathfrak{h}, \mathfrak{m}] \subset \mathfrak{m}$. For each X in the Lie algebra \mathfrak{g} of G , denote by X^* the Killing vector field on M generated by the one-parameter subgroup $\{\exp tX\}$ of G acting on M . The vector subspace \mathfrak{m} of $\hat{\mathfrak{g}} \subset \mathfrak{g}$ is identified with the tangent space $T_o(M)$ by the isomorphism $X \in \mathfrak{m} \rightarrow X_o^* \in T_o(M)$. The reductive decomposition $\hat{\mathfrak{g}} = \mathfrak{h} + \mathfrak{m}$ defines the homogeneous Riemannian structure $S = \nabla - \tilde{\nabla}$, where $\tilde{\nabla}$ is the canonical connection of $M = \hat{G}/H$ with respect to $\hat{\mathfrak{g}} = \mathfrak{h} + \mathfrak{m}$. The connection $\tilde{\nabla}$ is invariant under the action of \hat{G} , it is determined by $(\tilde{\nabla}_{X^*} Y^*)_o = [X^*, Y^*]_o = -[X, Y]_o^*$, for $X, Y \in \mathfrak{m}$, and the Ambrose-Singer equations (1.1) are satisfied. The tensor field S is also uniquely determined by its value at o because $M \equiv \hat{G}/H$ and S is \hat{G} -invariant. We have

$$(2.1) \quad (S_{X^*} Y^*)_o = (\nabla_{X^*} Y^*)_o + [X, Y]_o^* = \nabla_{Y_o^*} X^*, \quad X, Y \in \mathfrak{m}.$$

Moreover, as Ω is \hat{G} -invariant, from [15, Ch. 10, Prop. 2.7] it follows that $\tilde{\nabla}\Omega = 0$, so equations (1.2) are satisfied and S is also a homogeneous quaternionic Kähler structure.

Now, suppose that $(M = G/K, g, v)$ is a connected non-compact quaternion-Kähler symmetric space, and let $\hat{\mathfrak{g}} = \mathfrak{h} + \mathfrak{m}$ be a reductive decomposition corresponding to $M \equiv \hat{G}/H$ as above. We consider a Cartan decomposition $\mathfrak{g} = \mathfrak{k} + \mathfrak{p}$ of the Lie algebra \mathfrak{g} of G , and the Iwasawa decomposition $\mathfrak{g} = \mathfrak{k} + \mathfrak{a} + \mathfrak{n}$, where \mathfrak{k} is the Lie algebra of K , $\mathfrak{a} \subset \mathfrak{p}$ is a maximal \mathbb{R} -diagonalizable subalgebra of \mathfrak{g} , and \mathfrak{n} is a nilpotent subalgebra. Let A and N be the connected abelian and nilpotent Lie subgroups of G whose Lie algebras are \mathfrak{a} and \mathfrak{n} , respectively. The solvable Lie group AN acts simply transitively on M , so M is isometric to AN equipped with the left-invariant Riemannian metric defined by the scalar product induced on $\mathfrak{a} + \mathfrak{n} \cong \mathfrak{g}/\mathfrak{k} \cong \mathfrak{p}$ by a positive multiple of $B|_{\mathfrak{p} \times \mathfrak{p}}$, where B is the Killing form of \mathfrak{g} .

If $X \in \mathfrak{g} = \mathfrak{k} + \mathfrak{p}$, we write $X = X_{\mathfrak{k}} + X_{\mathfrak{p}}$, ($X_{\mathfrak{k}} \in \mathfrak{k}$, $X_{\mathfrak{p}} \in \mathfrak{p}$). We have the isomorphisms of vector spaces

$$(2.2) \quad \mathfrak{p} \cong \mathfrak{g}/\mathfrak{k} \cong \hat{\mathfrak{g}}/\mathfrak{h} \cong \mathfrak{m} \cong T_o(M) \cong \mathfrak{a} + \mathfrak{n},$$

with

$$\xi: \mathfrak{p} \xrightarrow{\cong} \mathfrak{m}, \quad \mu: \mathfrak{m} \xrightarrow{\cong} T_o(M), \quad \zeta: T_o(M) \xrightarrow{\cong} \mathfrak{a} + \mathfrak{n},$$

such that

$$\xi^{-1}(Z) = Z_{\mathfrak{p}}, \quad \mu(Z) = Z_o^*, \quad \zeta^{-1}(X) = X_o^*, \quad Z \in \mathfrak{m}, \quad X \in \mathfrak{a} + \mathfrak{n}.$$

For each $X \in \mathfrak{g}$, we have $(X_{\mathfrak{k}})_o^* = 0$ and $(\nabla(X_{\mathfrak{p}})^*)_o = 0$, and since the Levi-Civita connection ∇ has no torsion, for each $X, Y \in \mathfrak{g}$, we have

$$(2.3) \quad (\nabla_{X^*} Y^*)_o = (\nabla_{(X_{\mathfrak{p}})^*} (Y_{\mathfrak{k}})^*)_o = [(X_{\mathfrak{p}})^*, (Y_{\mathfrak{k}})^*]_o = -[X_{\mathfrak{p}}, Y_{\mathfrak{k}}]^*_o.$$

By (2.1) and (2.3), the homogeneous Riemannian structure S associated to a reductive decomposition $\hat{\mathfrak{g}} = \mathfrak{h} + \mathfrak{m}$ is given by

$$S_{X_o^*} Y_o^* = [X_{\mathfrak{k}}, Y_{\mathfrak{p}}]^*_o, \quad X, Y \in \mathfrak{m}.$$

Then, for each $X, Y \in \mathfrak{a} + \mathfrak{n}$, we have

$$(2.4) \quad S_{X_o^*} Y_o^* = S_{\xi(X_{\mathfrak{p}})_o^*} \xi(Y_{\mathfrak{p}})_o^* = [(\xi(X_{\mathfrak{p}}))_{\mathfrak{k}}, Y_{\mathfrak{p}}]^*_o.$$

This last formula will be helpful to give directly the associated homogeneous Riemannian structure S in terms of a basis of $\mathfrak{a} + \mathfrak{n}$. Moreover, we will get formula (2.6), which furnishes explicitly the coefficients θ^a , $a = 1, 2, 3$, in equations (1.3), for the quaternion-Hermitian structure on M which we will now describe. Before giving it, we note that formulas (2.4) and (2.6) will be useful in the next sections to find the homogeneous Riemannian structures on the spaces under study and their types as homogeneous quaternionic Kähler structures, respectively.

The quaternionic structure on M is defined by a 3-dimensional ideal $\mathfrak{u} \cong \mathfrak{sp}(1)$ of \mathfrak{k} , and we can suppose that $\mathfrak{u} = \langle E_1, E_2, E_3 \rangle$, with

$$[E_1, E_2] = 2E_3, \quad [E_2, E_3] = 2E_1, \quad [E_3, E_1] = 2E_2.$$

The next commutative diagram defines the complex structures $J_a \in \text{End}(\mathfrak{a} + \mathfrak{n})$ ($1 \leq a \leq 3$), which make $(\mathfrak{a} + \mathfrak{n}, \langle \cdot, \cdot \rangle, J_1, J_2, J_3)$ a quaternion-Hermitian vector space *isomorphic* to $(T_o(M), g_o, (J_1)_o, (J_2)_o, (J_3)_o)$.

$$(2.5) \quad \begin{array}{ccccccc} \mathfrak{p} & \xrightarrow{\xi} & \mathfrak{m} & \xrightarrow{\mu} & T_o(M) & \xrightarrow{\zeta} & \mathfrak{a} + \mathfrak{n} \\ \text{ad}_{E_a} \downarrow & & \downarrow & & (J_a)_o \downarrow & & \downarrow J_a \\ \mathfrak{p} & \xrightarrow{\xi} & \mathfrak{m} & \xrightarrow{\mu} & T_o(M) & \xrightarrow{\zeta} & \mathfrak{a} + \mathfrak{n} \end{array}$$

Now, the Lie subgroup of K generated by \mathfrak{u} is a normal subgroup isomorphic to $Sp(1)$, its isotropy representation is isomorphic to the representation of the normal subgroup $Sp(1)$ of $Sp(n)Sp(1)$ on \mathbb{R}^{4n} ($4n = \dim M$), and there exists a normal subgroup K_1 of K such that $K = K_1 Sp(1)$ and $K_1 \cap Sp(1)$ is contained in the center $\{\pm 1\}$ of $Sp(1)$. The Lie algebra \mathfrak{k}_1 of K_1 is an ideal of \mathfrak{k} such that $\mathfrak{k} = \mathfrak{u} \oplus \mathfrak{k}_1$, and we can get a basis $\mathcal{B} = \{E_1, E_2, E_3, \dots\}$ of \mathfrak{k} with the basic elements of \mathfrak{u} and some elements of \mathfrak{k}_1 . Let $\{\alpha_1, \alpha_2, \alpha_3, \dots\}$ be the dual basis of \mathcal{B} . Then we have

Proposition 2.4. *The homogeneous Riemannian structure S on the connected non-compact quaternion-Kähler symmetric space $M = G/K$ associated to the reductive decomposition $\hat{\mathfrak{g}} = \mathfrak{h} + \mathfrak{m}$ satisfies equations (1.3) with 1-forms θ^i , $i = 1, 2, 3$, given at $o \equiv H \in \hat{G}/H \equiv M$ by*

$$(2.6) \quad \theta^i(X_o^*) = 2\alpha_i((\xi(X_{\mathfrak{p}}))_{\mathfrak{k}}),$$

for each $X \in \mathfrak{a} + \mathfrak{n}$.

Proof. Let $X, Y \in \mathfrak{a} + \mathfrak{n}$. From the commutative diagram (2.5) we have $(J_a)_o Y_o^* = (\text{ad}_{E_a} Y_{\mathfrak{p}})_o^*$. Then, from (2.4), for each $a = 1, 2, 3$, we get

$$S_{X_o^*}((J_a)_o Y_o^*) = [(\xi(X_{\mathfrak{p}}))_{\mathfrak{k}}, \text{ad}_{E_a} Y_{\mathfrak{p}}]_o^*,$$

and

$$(J_a)_o(S_{X_o^*} Y_o^*) = (\text{ad}_{E_a} [(\xi(X_{\mathfrak{p}}))_{\mathfrak{k}}, Y_{\mathfrak{p}}]_o^*)^*.$$

Hence, applying Jacobi's identity, we deduce

$$(2.7) \quad S_{X_o^*}((J_a)_o Y_o^*) - (J_a)_o(S_{X_o^*} Y_o^*) = -[[E_a, (\xi(X_{\mathfrak{p}}))_{\mathfrak{k}}], Y_{\mathfrak{p}}]_o^*.$$

In terms of the basis \mathcal{B} , we have $(\xi(X_{\mathfrak{p}}))_{\mathfrak{k}} = \alpha_1((\xi(X_{\mathfrak{p}}))_{\mathfrak{k}})E_1 + \alpha_2((\xi(X_{\mathfrak{p}}))_{\mathfrak{k}})E_2 + \alpha_3((\xi(X_{\mathfrak{p}}))_{\mathfrak{k}})E_3 + \dots$, then

$$[E_1, (\xi(X_{\mathfrak{p}}))_{\mathfrak{k}}] = 2\alpha_2((\xi(X_{\mathfrak{p}}))_{\mathfrak{k}})E_3 - 2\alpha_3((\xi(X_{\mathfrak{p}}))_{\mathfrak{k}})E_2.$$

From (2.7), for instance for $a = 1$, we have

$$\begin{aligned} S_{X_o^*}((J_1)_o Y_o^*) - (J_1)_o(S_{X_o^*} Y_o^*) &= -[2\alpha_2((\xi(X_{\mathfrak{p}}))_{\mathfrak{k}})E_3 - 2\alpha_3((\xi(X_{\mathfrak{p}}))_{\mathfrak{k}})E_2, Y_{\mathfrak{p}}]_o^* \\ &= 2\alpha_3((\xi(X_{\mathfrak{p}}))_{\mathfrak{k}})[E_2, Y_{\mathfrak{p}}]_o^* - 2\alpha_2((\xi(X_{\mathfrak{p}}))_{\mathfrak{k}})[E_3, Y_{\mathfrak{p}}]_o^* \\ &= 2\alpha_3((\xi(X_{\mathfrak{p}}))_{\mathfrak{k}})(J_2)_o Y_o^* - 2\alpha_2((\xi(X_{\mathfrak{p}}))_{\mathfrak{k}})(J_3)_o Y_o^*, \end{aligned}$$

and equations (1.3) are satisfied, with $\theta^i(X_o^*)$, $i = 1, 2, 3$, as in the statement. \square

3 Homogeneous descriptions and structures

Since we want to obtain the types of homogeneous quaternionic Kähler structures on 8-dimensional non-compact quaternion-Kähler symmetric spaces, we need to get all the homogeneous descriptions of them.

As they are Alekseevsky spaces, we denote by $A_{SU(2,2)}$, $A_{Sp(2,1)}$ and $A_{G_2(2)}$, the spaces $SU(2,2)/S(U(2) \times U(2))$, $\mathbb{H}\mathbb{H}(2) = Sp(2,1)/(Sp(2) \times Sp(1))$, and $G_2(2)/SO(4)$, respectively. (For $\mathbb{H}\mathbb{H}(2)$ cf. [6, Ths. 1.1, 5.4]). Since the center of each of the corresponding full isometry groups is finite, Theorems 2.2 and 2.3 apply.

To find the corresponding homogeneous Riemannian structures, an important role will be played by the Lie algebra $\mathfrak{a} + \mathfrak{n}$ of the solvable factor AN in an Iwasawa decomposition of the full connected isometry group of each space.

To obtain each one of such structures S , we will use (2.4) and the identifications given by (2.2). In particular, we will also denote by $X \in \mathfrak{a} + \mathfrak{n}$ the vector $X_o^* = (X_p)_o^* \in T_o(M)$, and we will give the values $S_X Y = S_{X_o^*} Y_o^*$, for all X, Y in a suitable basis of $\mathfrak{a} + \mathfrak{n}$. We know that S is a homogeneous quaternionic Kähler structure and that Proposition 2.4 allows us to calculate directly the forms θ^a in (1.3).

3.1 The complex hyperbolic space $A_{SU(2,2)}$

3.1.1 The quaternion-Hermitian structure of $\mathfrak{a} + \mathfrak{n} \subset \mathfrak{su}(2, 2)$

The Lie algebra of $SU(2, 2)$ is

$$\mathfrak{su}(2, 2) = \left\{ \begin{pmatrix} A & B \\ \bar{B}^T & C \end{pmatrix} \in \mathfrak{sl}(4, \mathbb{C}) : A, C \in \mathfrak{u}(2) \right\}.$$

The involution τ of $\mathfrak{su}(2, 2)$ given by $\tau(X) = -\bar{X}^T$ defines the Cartan decomposition $\mathfrak{su}(2, 2) = \mathfrak{k} + \mathfrak{p}$, where $\mathfrak{k} = \mathfrak{s}(\mathfrak{u}(2) \oplus \mathfrak{u}(2))$. We consider the subspace \mathfrak{a} of \mathfrak{p} defined by

$$\mathfrak{a} = \left\{ \begin{pmatrix} 0 & 0 & 0 & s_2 \\ 0 & 0 & s_1 & 0 \\ 0 & s_1 & 0 & 0 \\ s_2 & 0 & 0 & 0 \end{pmatrix} : s_1, s_2 \in \mathbb{R} \right\},$$

which is a maximal \mathbb{R} -diagonalizable subalgebra of $\mathfrak{su}(2, 2)$. Let A_1 and A_2 be the generators of \mathfrak{a} defined by $(s_1, s_2) = (1, 0)$ and $(0, 1)$, respectively. Let f_1 and f_2 denote the linear functionals on \mathfrak{a} whose values on the above matrix of \mathfrak{a} are s_1 and s_2 , respectively. Then, the sets of positive roots and simple roots (with respect to a suitable order in \mathfrak{a}^*) are given by $\Sigma^+ = \{f_1 \pm f_2, 2f_1, 2f_2\}$ and $\Pi = \{f_1 - f_2, 2f_2\}$, respectively. The positive root vector spaces are

$$(3.1) \quad \mathfrak{g}_{f_1+f_2} = \left\{ \begin{pmatrix} 0 & z & -z & 0 \\ -\bar{z} & 0 & 0 & \bar{z} \\ -\bar{z} & 0 & 0 & \bar{z} \\ 0 & z & -z & 0 \end{pmatrix} \right\}, \quad \mathfrak{g}_{f_1-f_2} = \left\{ \begin{pmatrix} 0 & z & -z & 0 \\ -\bar{z} & 0 & 0 & -\bar{z} \\ -\bar{z} & 0 & 0 & -\bar{z} \\ 0 & -z & z & 0 \end{pmatrix} \right\},$$

$$(3.2) \quad \mathfrak{g}_{2f_1} = \left\{ \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & ix & -ix & 0 \\ 0 & ix & -ix & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \right\}, \quad \mathfrak{g}_{2f_2} = \left\{ \begin{pmatrix} ix & 0 & 0 & -ix \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ ix & 0 & 0 & -ix \end{pmatrix} \right\},$$

where $x \in \mathbb{R}$, $z \in \mathbb{C}$. The root vector spaces for the respective opposite roots are the corresponding sets of opposite conjugate transpose matrices. For each $f = f_1 \pm f_2$, let X_f and X'_f be the generators of \mathfrak{g}_f in (3.1) obtained by putting $z = 1$ and $z = i$, respectively, for each non-null entry. For $f = 2f_j$ ($1 \leq j \leq 2$), let U_f be the generator of \mathfrak{g}_f in (3.2) obtained by putting $x = 1$ in each non-zero entry. Let also X_f, X'_f, U_f be the corresponding elements of \mathfrak{g}_f for the respective opposite roots $f \in \Sigma \setminus \Sigma^+$. Now, we put $X_1 = X_{f_1+f_2}$, $Y_1 = X_{-f_1-f_2}$, $X'_1 = X'_{f_1+f_2}$, $Y'_1 = X'_{-f_1-f_2}$, $X_2 = X_{f_1-f_2}$, $Y_2 = X_{-f_1+f_2}$, $X'_2 = X'_{f_1-f_2}$, $Y'_2 = X'_{-f_1+f_2}$, $U_1 = U_{2f_1}$, $V_1 = U_{-2f_1}$, $U_2 = U_{2f_2}$, $V_2 = U_{-2f_2}$. Moreover, the centralizer of \mathfrak{a} in \mathfrak{k} is generated by $C = \text{diag}(i, -i, -i, i)$. Then,

$Z_{\mathfrak{k}}(\mathfrak{a}) + \mathfrak{a} = \langle C, A_1, A_2 \rangle$ is a Cartan subalgebra of $\mathfrak{su}(2, 2)$, and $\mathfrak{su}(2, 2) = (Z_{\mathfrak{k}}(\mathfrak{a}) + \mathfrak{a}) + \sum_{f \in \Sigma} \mathfrak{g}_f$ is the restricted-root space decomposition. We also have the Iwasawa decomposition $\mathfrak{su}(2, 2) = \mathfrak{k} + \mathfrak{a} + \mathfrak{n}$, where $\mathfrak{n} = \sum_{f \in \Sigma^+} \mathfrak{g}_f = \langle X_1, X'_1, X_2, X'_2, U_1, U_2 \rangle$.

The elements E_1, E_2, E_3 of $\mathfrak{k} = \mathfrak{s}(\mathfrak{u}(2) \oplus \mathfrak{u}(2)) \subset \mathfrak{su}(2, 2)$ given by

$$E_1 = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & -i & 0 \\ 0 & 0 & 0 & i \end{pmatrix}, \quad E_2 = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \end{pmatrix}, \quad E_3 = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & i \\ 0 & 0 & i & 0 \end{pmatrix},$$

satisfy $[E_1, E_2] = 2E_3$, $[E_2, E_3] = 2E_1$, $[E_3, E_1] = 2E_2$. The compact subalgebra $\mathfrak{u} \cong \mathfrak{sp}(1)$ generated by $\{E_1, E_2, E_3\}$ is an ideal of \mathfrak{k} , and the isotropy representation $\mathfrak{u} \rightarrow \mathfrak{gl}(\mathfrak{p})$ defines a quaternionic structure on $A_{SU(2,2)}$. From the action of each E_i on \mathfrak{p} and through the isomorphisms $\mathfrak{p} \cong \mathfrak{su}(2, 2)/\mathfrak{k} \cong \mathfrak{a} + \mathfrak{n}$, we obtain the complex structures J_a ($a = 1, 2, 3$), acting on $\mathfrak{a} + \mathfrak{n}$. The action on the elements A_j, U_j, X_j, X'_j ($j = 1, 2$) of the basis of $\mathfrak{a} + \mathfrak{n}$ is the following.

	A_1	A_2	U_1	U_2	X_1	X'_1	X_2	X'_2
J_1	$-U_1$	U_2	A_1	$-A_2$	X'_1	$-X_1$	X'_2	$-X_2$
J_2	$\frac{1}{2}(X_1 - X_2)$	$\frac{1}{2}(X_1 + X_2)$	$\frac{1}{2}(X'_1 - X'_2)$	$-\frac{1}{2}(X'_1 + X'_2)$	$-A_1 - A_2$	$-U_1 + U_2$	$A_1 - A_2$	$U_1 + U_2$
J_3	$\frac{1}{2}(X'_1 - X'_2)$	$\frac{1}{2}(X'_1 + X'_2)$	$-\frac{1}{2}(X_1 - X_2)$	$\frac{1}{2}(X_1 + X_2)$	$U_1 - U_2$	$-A_1 - A_2$	$-U_1 - U_2$	$A_1 - A_2$

Any positive multiple of the restriction of the Killing form B to $\mathfrak{p} \times \mathfrak{p}$ defines a Hermitian metric adapted to (J_1, J_2, J_3) . We consider the scalar product $\langle \cdot, \cdot \rangle$ induced in $\mathfrak{a} + \mathfrak{n}$ by the isomorphism $\mathfrak{p} \cong \mathfrak{a} + \mathfrak{n}$ and $\frac{1}{16}B|_{\mathfrak{p} \times \mathfrak{p}}$. This product makes the basis orthogonal, with $\langle A_j, A_j \rangle = \langle U_j, U_j \rangle = 1$, $\langle X_j, X_j \rangle = \langle X'_j, X'_j \rangle = 2$, and $(\mathfrak{a} + \mathfrak{n}, \langle \cdot, \cdot \rangle, J_1, J_2, J_3)$ is a quaternion-Hermitian vector space.

3.1.2 Homogeneous descriptions and structures of $A_{SU(2,2)}$

We now determine the homogeneous descriptions and structures of $A_{SU(2,2)}$. These descriptions will be obtained from the parabolic subalgebras of $\mathfrak{su}(2, 2)$ and their refined Langland decompositions, by using Theorems 2.2 and 2.3. The four standard parabolic subalgebras are parametrized by the subsets $\Pi, \emptyset, \Psi_1 = \{f_1 - f_2\}$ and $\Psi_2 = \{2f_2\}$ of Π .

The case $\Psi = \Pi$. We have $\mathfrak{e}'_{\Pi} = \mathfrak{a}_{\Pi} = \mathfrak{n}_{\Pi} = \{0\}$, hence the refined Langlands decomposition is $\mathfrak{p}_{\Pi} = \mathfrak{su}(2, 2) + \{0\} + \{0\} + \{0\}$. Thus, for a connected closed co-compact subgroup \hat{G} of P_{Π} , it is only possible to have $\hat{G} = L = L'_{\Pi} = SU(2, 2)$. Then the only transitive action on $A_{SU(2,2)}$ coming from $\Psi = \Pi$ is that of the full isometry group $SU(2, 2)$, and we have the description of $A_{SU(2,2)}$ as the symmetric space $SU(2, 2)/S(U(2) \times U(2))$. The associated reductive decomposition is the Cartan decomposition $\mathfrak{su}(2, 2) = \mathfrak{s}(\mathfrak{u}(2) \oplus \mathfrak{u}(2)) + \mathfrak{p}$, and the corresponding homogeneous quaternionic Kähler structure is $S = 0$.

The case $\Psi = \emptyset$. In this case, $\mathfrak{l}' = \mathfrak{l}'_{\emptyset} + \mathfrak{e}'_{\emptyset} + \mathfrak{a}_{\emptyset}$ is commutative, $\mathfrak{l}'_{\emptyset} = \{0\}$, $\mathfrak{e}'_{\emptyset} = Z_{\mathfrak{k}}(\mathfrak{a})$, and $\mathfrak{a}_{\emptyset} = \mathfrak{a}$. Hence the refined Langlands decomposition of the minimal parabolic subalgebra is $\mathfrak{p}_{\emptyset} = \{0\} + Z_{\mathfrak{k}}(\mathfrak{a}) + \mathfrak{a} + \mathfrak{n}$. For each connected

closed subgroup E of $E'_0 A \cong U(1) \times \mathbb{R}^2$ we get a co-compact subgroup EN of $SU(2, 2)$, where N is the nilpotent factor in the Iwasawa decomposition of $SU(2, 2)$. In order to obtain a transitive action it is sufficient that the projection of the Lie algebra $\mathfrak{e} \subset Z_{\mathfrak{k}}(\mathfrak{a}) + \mathfrak{a}$ of E to \mathfrak{a} be surjective.

For each choice of a two-dimensional subspace \mathfrak{e} of $Z_{\mathfrak{k}}(\mathfrak{a}) + \mathfrak{a} = \langle C, A_1, A_2 \rangle$ such as $C \notin \mathfrak{e}$, we obtain a Lie subalgebra $\mathfrak{e} + \mathfrak{n}$ of $\mathfrak{su}(2, 2)$ which generates a connected closed subgroup $\hat{G} = EN$ which acts simply transitively on $A_{SU(2,2)}$. We can suppose that $\mathfrak{e} = \mathfrak{e}_{\lambda, \mu}$ is generated by two elements of the form $\lambda C + A_1$, $\mu C + A_2$, with $\lambda, \mu \in \mathbb{R}$. This gives the reductive decomposition $\hat{\mathfrak{g}}^{\lambda, \mu} = \{0\} + \hat{\mathfrak{g}}^{\lambda, \mu}$ associated to the description $A_{SU(2,2)} = E_{\lambda, \mu} N$, where $\hat{\mathfrak{g}}^{\lambda, \mu} = \langle \lambda C + A_1, \mu C + A_2, U_1, U_2, X_1, X'_1, X_2, X'_2 \rangle$. Distinct values of λ, μ define non-isomorphic Lie algebras. If $\lambda = \mu = 0$, we have the choice $\mathfrak{e} = \mathfrak{a}$, which gives the usual description of $A_{SU(2,2)}$ as the solvable Lie group AN .

We get a two-parameter family of structures $S^{\lambda, \mu}$ associated to the family of reductive decompositions $\hat{\mathfrak{g}}^{\lambda, \mu} = \{0\} + \hat{\mathfrak{g}}^{\lambda, \mu}$. With the identifications in (2.2), where now $\mathfrak{m} = \hat{\mathfrak{g}}^{\lambda, \mu} = \mathfrak{e}_{\lambda, \mu} + \mathfrak{n}$, and by using (2.4), we obtain that $S = S^{\lambda, \mu}$ is given at o by Table I, with $\delta = \varepsilon = 1$.

Table I	A_1	A_2	U_1	U_2	X_1	X'_1	X_2	X'_2
S_{A_1}	0	0	0	0	$2\lambda X'_1$	$-2\lambda X_1$	$2\lambda X'_2$	$-2\lambda X_2$
S_{A_2}	0	0	0	0	$2\mu X'_1$	$-2\mu X_1$	$2\mu X'_2$	$-2\mu X_2$
S_{U_1}	$-2U_1$	0	$2A_1$	0	X'_2	$-X_2$	X'_1	$-X_1$
S_{U_2}	0	$-2\delta U_2$	0	$2\delta A_2$	$\delta X'_2$	$-\delta X_2$	$\delta X'_1$	$-\delta X_1$
S_{X_1}	$-X_1$	$-X_1$	X'_2	X'_2	$2(A_1 + A_2)$	0	0	$-2(U_1 + U_2)$
$S_{X'_1}$	$-X'_1$	$-X'_1$	$-X_2$	$-X_2$	0	$2(A_1 + A_2)$	$2(U_1 + U_2)$	0
S_{X_2}	$-\varepsilon X_2$	εX_2	$\varepsilon X'_1$	$-\varepsilon X'_1$	0	$-2\varepsilon(U_1 - U_2)$	$2\varepsilon(A_1 - A_2)$	0
$S_{X'_2}$	$-\varepsilon X'_2$	$\varepsilon X'_2$	$-\varepsilon X_1$	εX_1	$2\varepsilon(U_1 - U_2)$	0	0	$2\varepsilon(A_1 - A_2)$

Each structure $S^{\lambda, \mu}$ is also characterized by the fact that $\tilde{\nabla} = \nabla - S^{\lambda, \mu}$ is the canonical connection for the Lie group $E_{\lambda, \mu} N$, which is the connection for which every left-invariant vector field on $E_{\lambda, \mu} N$ is parallel.

The choice $\mathfrak{e} = Z_{\mathfrak{k}}(\mathfrak{a}) + \mathfrak{a}$ gives $A_{SU(2,2)} \equiv (U(1) \times \mathbb{R}^2)N/U(1)$, and the associated reductive decomposition is $\hat{\mathfrak{g}} = \mathfrak{h} + \mathfrak{m}$, where $\hat{\mathfrak{g}} = \mathfrak{p}_0$, $\mathfrak{h} = Z_{\mathfrak{k}}(\mathfrak{a})$ and $\mathfrak{m} = \mathfrak{a} + \mathfrak{n}$. The corresponding structure coincides with the above structure $S^{\lambda, \mu}$, for $\lambda = \mu = 0$.

The case $\Psi = \Psi_1$. Then $[\Psi_1] = \{\pm(f_1 - f_2)\}$, and one has $\mathfrak{l}'_{\Psi_1} = \langle A_2 - A_1 \rangle + Z_{\mathfrak{k}}(\mathfrak{a}) + \langle X_2, Y_2, X'_2, Y'_2 \rangle$, $\mathfrak{e}'_{\Psi_1} = \{0\}$, $\mathfrak{a}_{\Psi_1} = \langle A_1 + A_2 \rangle$, and $\mathfrak{n}_{\Psi_1} = \mathfrak{g}_{f_1+f_2} + \mathfrak{g}_{2f_1} + \mathfrak{g}_{2f_2} = \langle X_1, X'_1, U_1, U_2 \rangle$. Hence, the corresponding parabolic subalgebra is $\mathfrak{p}_{\Psi_1} = \mathfrak{l}'_{\Psi_1} + \{0\} + \mathfrak{a}_{\Psi_1} + \mathfrak{n}_{\Psi_1}$, with

$$\mathfrak{l}'_{\Psi_1} = \left\{ \begin{pmatrix} ir & z & w & s \\ -\bar{z} & -ir & -s & \bar{w} \\ \bar{w} & -s & -ir & -\bar{z} \\ s & w & z & ir \end{pmatrix} : r, s \in \mathbb{R}, z, w \in \mathbb{C} \right\} \cong \mathfrak{sl}(2, \mathbb{C}).$$

The connected subgroup $E = A_{\Psi_1} \cong \mathbb{R}$ of P_{Ψ_1} with Lie algebra \mathfrak{a}_{Ψ_1} is the only possible choice to get a co-compact subgroup $\hat{G} = LEN_{\Psi_1}$ of $SU(2, 2)$, with $L = L'_{\Psi_1} \cong Sl(2, \mathbb{C})$, which acts transitively on $A_{SU(2,2)}$. The isotropy algebra is $\mathfrak{h} = \hat{\mathfrak{g}} \cap \mathfrak{k} = \mathfrak{l}'_{\Psi_1} \cap \mathfrak{s}(\mathfrak{u}(2) \oplus \mathfrak{u}(2)) = \langle C, (X_2)_{\mathfrak{k}}, (X'_2)_{\mathfrak{k}} \rangle \cong \mathfrak{su}(2)$, and

we get the reductive decomposition $\hat{\mathfrak{g}} = \mathfrak{h} + \mathfrak{m}$, where $\hat{\mathfrak{g}} = \mathfrak{p}_{\Psi_1}$, and $\mathfrak{m} = \langle A_1, A_2, U_1, U_2, X_1, X'_1, (X_2)_{\mathfrak{p}}, (X'_2)_{\mathfrak{p}} \rangle$, which is associated to the description $A_{SU(2,2)} \equiv Sl(2, \mathbb{C})\mathbb{R}N_{\Psi_1}/SU(2)$. By using (2.4), we obtain that the corresponding structure S is given at o by Table I, with $\lambda = \mu = \varepsilon = 0$, $\delta = 1$.

The case $\Psi = \Psi_2$. Then $[\Psi_2] = \{\pm 2f_2\}$, and we have $\mathfrak{l}'_{\Psi_2} = \langle A_2, U_2, V_2 \rangle$, $\mathfrak{e}'_{\Psi_2} = Z_{\mathfrak{k}}(\mathfrak{a})$, $\mathfrak{a}_{\Psi_2} = \langle A_1 \rangle$, $\mathfrak{n}_{\Psi_2} = \mathfrak{g}_{f_1+f_2} + \mathfrak{g}_{f_1-f_2} + \mathfrak{g}_{2f_1} = \langle X_1, X'_1, X_2, X'_2, U_1 \rangle$. The corresponding parabolic subalgebra is $\mathfrak{p}_{\Psi_2} = \mathfrak{l}'_{\Psi_2} + Z_{\mathfrak{k}}(\mathfrak{a}) + \mathfrak{a}_{\Psi_2} + \mathfrak{n}_{\Psi_2}$, with

$$\mathfrak{l}'_{\Psi_2} = \left\{ \begin{pmatrix} ix & 0 & 0 & z \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ \bar{z} & 0 & 0 & -ix \end{pmatrix} : x \in \mathbb{R}, z \in \mathbb{C} \right\} \cong \mathfrak{su}(1, 1) \cong \mathfrak{sl}(2, \mathbb{R}).$$

For each connected closed subgroup E of P_{Ψ_2} whose Lie algebra \mathfrak{e} is a non-trivial subspace of $\mathfrak{e}'_{\Psi_2} + \mathfrak{a}_{\Psi_2}$, $\mathfrak{e} \neq \mathfrak{e}'_{\Psi_2}$, we get a co-compact subgroup $\hat{G} = LEN_{\Psi_2}$ of $SU(2, 2)$, with $L = L'_{\Psi_2} \cong SU(1, 1) \cong Sl(2, \mathbb{R})$, which acts transitively on $A_{SU(2,2)}$.

If $\dim E = 1$, we get the description $A_{SU(2,2)} \equiv SU(1, 1)\mathbb{R}N_{\Psi_2}/U(1)$. In fact, if \mathfrak{e} is an one-dimensional subspace of $\mathfrak{e}'_{\Psi_2} + \mathfrak{a}_{\Psi_2}$, with $\mathfrak{e} \neq \mathfrak{e}'_{\Psi_2}$, then $\mathfrak{e} = \mathfrak{e}_{\lambda}$ is spanned by an element of the form $\lambda C + A_1$, for $\lambda \in \mathbb{R}$, so one gets a one-parameter family of reductive decompositions $\hat{\mathfrak{g}}^{\lambda} = \mathfrak{h} + \mathfrak{m}^{\lambda}$, where $\mathfrak{h} = \hat{\mathfrak{g}} \cap \mathfrak{k} = \mathfrak{l}'_{\Psi_2} \cap \mathfrak{s}(\mathfrak{u}(2) \oplus \mathfrak{u}(2)) = \langle (U_2)_{\mathfrak{k}} \rangle \cong \mathfrak{u}(1)$ is the isotropy algebra and $\mathfrak{m}^{\lambda} = \langle \lambda C + A_1, A_2, U_1, (U_2)_{\mathfrak{p}}, X_1, X'_1, X_2, X'_2 \rangle$. By (2.4), we have that the corresponding one-parameter family of structures S^{λ} is given at o by Table I, with $\mu = \delta = 0$, $\varepsilon = 1$. In particular, $\lambda = 0$ corresponds to the choice $\mathfrak{e} = \mathfrak{a}_{\Psi_2}$.

If $\dim E = 2$, then $E \cong U(1) \times \mathbb{R}$, which gives $A_{SU(2,2)} \equiv SU(1, 1)(U(1) \times \mathbb{R})N_{\Psi_2}/(U(1) \times U(1))$. The natural associated reductive decomposition is $\hat{\mathfrak{g}}' = \mathfrak{h}' + \mathfrak{m}'$, where $\hat{\mathfrak{g}} = \mathfrak{p}_{\Psi_2}$, $\mathfrak{h}' = \langle C, (U_2)_{\mathfrak{k}} \rangle \cong \mathfrak{u}(1) \oplus \mathfrak{u}(1)$, $\mathfrak{m}' = \langle A_1, A_2, U_1, (U_2)_{\mathfrak{p}}, X_1, X'_1, X_2, X'_2 \rangle$, and the corresponding structure S' coincides with the previous S^{λ} , for $\lambda = 0$.

In particular, we have proved

Proposition 3.1. *The space $A_{SU(2,2)}$ admits the homogeneous quaternionic Kähler structures S (associated to the corresponding homogeneous description \hat{G}/H) given by*

S	\hat{G}/H
0	$SU(2, 2)/S(U(2) \times U(2))$
Table I ($\lambda, \mu \in \mathbb{R}$; $\delta = \varepsilon = 1$)	$E_{\lambda, \mu}N$ ($E_{\lambda, \mu} \cong \mathbb{R}^2$)
Table I ($\lambda = \mu = \varepsilon = 0$, $\delta = 1$)	$Sl(2, \mathbb{C})\mathbb{R}N_{\Psi_1}/SU(2)$
Table I ($\lambda \in \mathbb{R}$; $\mu = \delta = 0$, $\varepsilon = 1$)	$SU(1, 1)\mathbb{R}N_{\Psi_2}/U(1)$

where N is the nilpotent factor in the Iwasawa decomposition $SU(2, 2) = S(U(2) \times U(2))AN$ and Ψ_1 and Ψ_2 are the subsets of simple restricted roots defining the intermediate parabolic subalgebras of $(\mathfrak{su}(2, 2), \mathfrak{a})$. \square

3.2 The quaternionic hyperbolic space $A_{Sp(2,1)} = \mathbb{H}\mathbb{H}(2)$

3.2.1 The quaternion-Hermitian structure of $\mathfrak{a} + \mathfrak{n} \subset \mathfrak{sp}(2, 1)$

The Lie algebra $\mathfrak{sp}(2, 1)$ can be described as the subalgebra of $\mathfrak{gl}(6, \mathbb{C})$ of all matrices $X \in \mathfrak{sp}(3, \mathbb{C})$ satisfying $\bar{X}^T J_{2,1} + J_{2,1} X = 0$, where $J_{2,1} = \text{diag}(-1, -1, 1, -1, -1, 1)$. The involution τ of $\mathfrak{sp}(2, 1)$ given by $\tau(X) = -\bar{X}^T$ defines the Cartan decomposition $\mathfrak{sp}(2, 1) = \mathfrak{k} + \mathfrak{p}$, where

$$\mathfrak{k} = \left\{ \begin{pmatrix} ia & z & 0 & u & w & 0 \\ -\bar{z} & ib & 0 & w & v & 0 \\ 0 & 0 & ic & 0 & 0 & \alpha \\ -\bar{u} & -\bar{w} & 0 & -ia & \bar{z} & 0 \\ -\bar{v} & -\bar{v} & 0 & -z & -ib & 0 \\ 0 & 0 & -\bar{\alpha} & 0 & 0 & -ic \end{pmatrix} : \begin{matrix} a, b, c \in \mathbb{R}, \\ z, u, v, w, \alpha \in \mathbb{C} \end{matrix} \right\} \cong \mathfrak{sp}(2) \oplus \mathfrak{sp}(1),$$

$$\mathfrak{p} = \left\{ \begin{pmatrix} 0 & 0 & p_1 & 0 & 0 & q_1 \\ 0 & 0 & p_2 & 0 & 0 & q_2 \\ \bar{p}_1 & \bar{p}_2 & 0 & q_1 & q_2 & 0 \\ 0 & 0 & \bar{q}_1 & 0 & 0 & -\bar{p}_1 \\ 0 & 0 & \bar{q}_2 & 0 & 0 & -\bar{p}_2 \\ \bar{q}_1 & \bar{q}_2 & 0 & -p_1 & -p_2 & 0 \end{pmatrix} : p_1, p_2, q_1, q_2 \in \mathbb{C} \right\}.$$

The element A_0 of \mathfrak{p} obtained by taking $p_1 = 1$ and $p_2 = q_1 = q_2 = 0$ generates a maximal \mathbb{R} -diagonalizable subalgebra \mathfrak{a} of $\mathfrak{sp}(2, 1)$. The set of roots corresponding to \mathfrak{a} is $\Sigma = \{\pm f_0, \pm 2f_0\}$, where $f_0 \in \mathfrak{a}^*$ is given by $f_0(A_0) = 1$. The set $\Pi = \{f_0\}$ is a system of simple roots and the corresponding positive root system is $\Sigma^+ = \{f_0, 2f_0\}$. The positive root vector spaces are

$$(3.3) \quad \mathfrak{g}_{f_0} = \left\{ \begin{pmatrix} 0 & z & 0 & 0 & w & 0 \\ -\bar{z} & 0 & \bar{z} & w & 0 & w \\ 0 & z & 0 & 0 & w & 0 \\ 0 & -\bar{w} & 0 & 0 & \bar{z} & 0 \\ -\bar{w} & 0 & \bar{w} & -z & 0 & -z \\ 0 & \bar{w} & 0 & 0 & -\bar{z} & 0 \end{pmatrix} : z, w \in \mathbb{C} \right\},$$

$$(3.4) \quad \mathfrak{g}_{2f_0} = \left\{ \begin{pmatrix} ia & 0 & -ia & u & 0 & u \\ 0 & 0 & 0 & 0 & 0 & 0 \\ ia & 0 & -ia & u & 0 & u \\ -\bar{u} & 0 & \bar{u} & -ia & 0 & -ia \\ 0 & 0 & 0 & 0 & 0 & 0 \\ \bar{u} & 0 & -\bar{u} & ia & 0 & ia \end{pmatrix} : a \in \mathbb{R}, u \in \mathbb{C} \right\}.$$

The root vector spaces for the respective opposite roots are the corresponding sets of opposite conjugate transpose matrices. Now, let X_1, U_1 and U'_1 be the elements of \mathfrak{g}_{2f_0} in (3.4) obtained by putting $(a, u) = (-1, 0), (0, 1)$ and $(0, i)$, respectively, and denote by Y_1, V_1 and V'_1 the corresponding respective elements of \mathfrak{g}_{-2f_0} . On the other hand, let X_2, X'_2, U_2 and U'_2 be the elements of \mathfrak{g}_{f_0} in (3.3) obtained by setting $(z, w) = (1, 0), (-i, 0), (0, 1)$ and $(0, i)$, respectively, and let Y_2, Y'_2, V_2 and V'_2 denote the corresponding respective elements of \mathfrak{g}_{-f_0} . Then, $\mathfrak{g}_{f_0} = \langle X_2, X'_2, U_2, U'_2 \rangle$, $\mathfrak{g}_{2f_0} = \langle X_1, U_1, U'_1 \rangle$, $\mathfrak{g}_{-f_0} = \langle Y_2, Y'_2, V_2, V'_2 \rangle$, $\mathfrak{g}_{-2f_0} = \langle Y_1, V_1, V'_1 \rangle$, and we have the Iwasawa decomposition $\mathfrak{sp}(2, 1) = \mathfrak{k} + \mathfrak{a} + \mathfrak{n}$, where

$\mathfrak{n} = \mathfrak{g}_{f_0} + \mathfrak{g}_{2f_0} = \langle X_1, U_1, U'_1, X_2, X'_2, U_2, U'_2 \rangle$. The centralizer of \mathfrak{a} in \mathfrak{k} is

$$(3.5) \quad Z_{\mathfrak{k}}(\mathfrak{a}) = \left\{ \begin{pmatrix} ia & 0 & 0 & u & 0 & 0 \\ 0 & ib & 0 & 0 & v & 0 \\ 0 & 0 & ia & 0 & 0 & -u \\ -\bar{u} & 0 & 0 & -ia & 0 & 0 \\ 0 & -\bar{v} & 0 & 0 & -ib & 0 \\ 0 & 0 & \bar{u} & 0 & 0 & -ia \end{pmatrix} : \begin{matrix} a, b \in \mathbb{R}, \\ u, v \in \mathbb{C} \end{matrix} \right\} \cong \mathfrak{sp}(1) \oplus \mathfrak{sp}(1),$$

and it is generated by the elements C_1, C_2, C_3, D_1, D_2 and D_3 in (3.5) obtained by taking $(a, b, u, v) = (1, 0, 0, 0), (0, 0, 1, 0), (0, 0, i, 0), (0, 1, 0, 0), (0, 0, 0, 1)$ and $(0, 0, 0, i)$, respectively. The subspaces $\langle C_1, C_2, C_3 \rangle$ and $\langle D_1, D_2, D_3 \rangle$ are ideals of $Z_{\mathfrak{k}}(\mathfrak{a})$ isomorphic to $\mathfrak{sp}(1)$.

The elements E_1, E_2, E_3 of $\mathfrak{k} = \mathfrak{sp}(2) \oplus \mathfrak{sp}(1) \subset \mathfrak{sp}(2, 1)$ given by

$$\begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & i & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -i \end{pmatrix}, \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}, \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & i \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & i & 0 & 0 & 0 \end{pmatrix},$$

respectively, satisfy $[E_1, E_2] = 2E_3$, $[E_2, E_3] = 2E_1$, $[E_3, E_1] = 2E_2$, and generate a compact ideal $\mathfrak{u} \cong \mathfrak{sp}(1)$ of \mathfrak{k} . The isotropy representation $\mathfrak{u} \rightarrow \mathfrak{gl}(\mathfrak{p})$ defines a quaternionic structure on $A_{Sp(2,1)}$. The isomorphisms $\mathfrak{p} \cong \mathfrak{sp}(2, 1)/\mathfrak{k} \cong \mathfrak{a} + \mathfrak{n}$ allow to obtain the complex structures J_a ($a = 1, 2, 3$), acting on $\mathfrak{a} + \mathfrak{n}$. The action on the elements of the basis $\{A_0, X_1, U_1, U'_1, X_2, X'_2, U_2, U'_2\}$ of $\mathfrak{a} + \mathfrak{n}$ is given in the following table.

	A_0	X_1	U_1	U'_1	X_2	X'_2	U_2	U'_2
J_1	$-X_1$	A_0	U'_1	$-U_1$	$-X'_2$	X_2	U'_2	$-U_2$
J_2	$-U_1$	$-U'_1$	A_0	X_1	$-U_2$	$-U'_2$	X_2	X'_2
J_3	$-U'_1$	U_1	$-X_1$	A_0	$-U'_2$	U_2	$-X'_2$	X_2

The above basis of $\mathfrak{a} + \mathfrak{n}$ is orthonormal with respect to the scalar product $\langle \cdot, \cdot \rangle$ induced in $\mathfrak{a} + \mathfrak{n}$ by the isomorphism $\mathfrak{p} \cong \mathfrak{a} + \mathfrak{n}$ and $\frac{1}{32}B|_{\mathfrak{p} \times \mathfrak{p}}$, where B is the Killing form of $\mathfrak{sp}(2, 1)$, and $(\mathfrak{a} + \mathfrak{n}, \langle \cdot, \cdot \rangle, J_1, J_2, J_3)$ is a quaternion-Hermitian vector space.

3.2.2 Homogeneous descriptions and structures of $A_{Sp(2,1)}$

We now determine the homogeneous descriptions and structures of $A_{Sp(2,1)}$. The two only parabolic subalgebras of $\mathfrak{sp}(2, 1)$ are parametrized by the subsets Π and \emptyset of $\Pi = \{f_0\}$. We will use them and their refined Langland decompositions to obtain, by means of Theorems 2.2 and 2.3, the homogeneous descriptions of $A_{Sp(2,1)} = \mathbb{H}\mathbb{H}(2)$.

The case $\Psi = \Pi$. We have $\mathfrak{e}'_{\Pi} = \mathfrak{a}_{\Pi} = \mathfrak{n}_{\Pi} = \{0\}$, hence the refined Langlands decomposition is $\mathfrak{p}_{\Pi} = \mathfrak{sp}(2, 1) + \{0\} + \{0\} + \{0\}$. Theorem 2.2 then says that the only connected closed co-compact subgroup of P_{Π} is $\hat{G} = Sp(2, 1)$. Thus the only transitive action coming from $\Psi = \Pi$ is that of the full isometry group $Sp(2, 1)$. This gives the description of $A_{Sp(2,1)} = \mathbb{H}\mathbb{H}(2)$ as the symmetric space

$Sp(2, 1)/(Sp(2) \times Sp(1))$. The associated reductive decomposition is the Cartan decomposition $\mathfrak{sp}(2, 1) = (\mathfrak{sp}(2) \oplus \mathfrak{sp}(1)) + \mathfrak{p}$, and the corresponding structure is $S = 0$.

The case $\Psi = \emptyset$. We have $\mathfrak{l}' = \mathfrak{a} + Z_{\mathfrak{k}}(\mathfrak{a}) = \mathfrak{l}'_{\emptyset} + \mathfrak{e}'_{\emptyset} + \mathfrak{a}_{\emptyset}$, with $\mathfrak{l}'_{\emptyset} = \{0\}$, $\mathfrak{e}'_{\emptyset} = Z_{\mathfrak{k}}(\mathfrak{a})$, $\mathfrak{a}_{\emptyset} = \mathfrak{a}$, so the refined Langlands decomposition of the corresponding parabolic subalgebra is $\mathfrak{p}_{\emptyset} = \{0\} + Z_{\mathfrak{k}}(\mathfrak{a}) + \mathfrak{a} + \mathfrak{n} = \{0\} + (\mathfrak{sp}(1) \oplus \mathfrak{sp}(1)) + \mathfrak{a} + (\mathfrak{g}_{f_0} + \mathfrak{g}_{2f_0})$. For each connected closed subgroup E of $E'_{\emptyset}A \cong Sp(1)Sp(1)\mathbb{R}$ we get a co-compact subgroup EN of $Sp(2, 1)$. In order to get a transitive action on $A_{Sp(2, 1)}$ it is sufficient, by Theorem 2.3, that the projection of the Lie algebra $\mathfrak{e} \subset Z_{\mathfrak{k}}(\mathfrak{a}) + \mathfrak{a}$ to \mathfrak{a} be surjective.

Now, for each one-dimensional subspace \mathfrak{e} of $Z_{\mathfrak{k}}(\mathfrak{a}) + \mathfrak{a} = \langle C_j, D_j, A_0 \rangle_{j=1,2,3}$ such that the projection of \mathfrak{e} to \mathfrak{a} is an isomorphism, we get a Lie subalgebra $\mathfrak{e} + \mathfrak{n}$ of $\mathfrak{sp}(2, 1)$ which generates a connected closed Lie subgroup $\hat{G} = EN$ acting simply transitively on $A_{Sp(2, 1)}$. We can suppose that $\mathfrak{e} = \mathfrak{e}_{\lambda, \mu}$ is generated by one element of the form $\hat{A}_0 = A_0 + \sum_{j=1}^3 (\lambda_j C_j + \mu_j D_j)$, with $\lambda = (\lambda_1, \lambda_2, \lambda_3)$, $\mu = (\mu_1, \mu_2, \mu_3) \in \mathbb{R}^3$. This gives the reductive decomposition $\hat{\mathfrak{g}}^{\lambda, \mu} = \{0\} + \hat{\mathfrak{g}}^{\lambda, \mu}$ associated to the description $A_{Sp(2, 1)} = E_{\lambda, \mu}N$, where $\hat{\mathfrak{g}}^{\lambda, \mu} = \langle \hat{A}_0, X_1, U_1, U'_1, X_2, X'_2, U_2, U'_2 \rangle$. If $\lambda = \mu = (0, 0, 0)$, we have $\mathfrak{e} = \mathfrak{a}$, which gives the usual description of $A_{Sp(2, 1)}$ as the solvable Lie group AN . Then we have a six-parameter family of structures $S = S^{\lambda, \mu}$ associated to the reductive decompositions $\hat{\mathfrak{g}}^{\lambda, \mu} = \{0\} + \hat{\mathfrak{g}}^{\lambda, \mu}$. For each $\lambda, \mu \in \mathbb{R}^3$, and with the identifications in (2.2), where now $\mathfrak{m} = \hat{\mathfrak{g}}^{\lambda, \mu} = \mathfrak{e}_{\lambda, \mu} + \mathfrak{n}$, we apply (2.4) to obtain the homogeneous quaternionic Kähler structure $S = S^{\lambda, \mu}$, which is given at o by Table II.

Table II	A_0	X_1	U_1	U'_1	X_2	X'_2	U_2	U'_2
S_{A_0}	0	$2\lambda_2 U'_1$ $-2\lambda_3 U_1$	$2\lambda_1 U'_1$ $+2\lambda_3 X_1$	$-2\lambda_1 U_1$ $-2\lambda_2 X_1$	$(\mu_1 - \lambda_1)X'_2$ $+(\lambda_2 - \mu_2)U'_2$ $+(\lambda_3 - \mu_3)U'_2$	$(\lambda_1 - \mu_1)X_2$ $+(\lambda_2 + \mu_2)U'_2$ $-(\lambda_3 + \mu_3)U_2$	$(\lambda_1 + \mu_1)U'_2$ $+(\mu_2 - \lambda_2)X_2$ $+(\lambda_3 + \mu_3)X'_2$	$-(\lambda_1 + \mu_1)U_2$ $-(\lambda_2 + \mu_2)X'_2$ $+(\mu_3 - \lambda_3)X_2$
S_{X_1}	$-2X_1$	$2A_0$	0	0	$-X'_2$	X_2	U'_2	$-U_2$
S_{U_1}	$-2U_1$	0	$2A_0$	0	$-U'_2$	$-U'_2$	X_2	X'_2
$S_{U'_1}$	$-2U'_1$	0	0	$2A_0$	$-U'_2$	U_2	$-X'_2$	X_2
S_{X_2}	$-X_2$	$-X'_2$	$-U_2$	$-U'_2$	A_0	X_1	U_1	U'_1
$S_{X'_2}$	$-X'_2$	X_2	$-U'_2$	U_2	$-X_1$	A_0	$-U'_1$	U_1
S_{U_2}	$-U_2$	U'_2	X_2	$-X'_2$	$-U_1$	U'_1	A_0	$-X_1$
$S_{U'_2}$	$-U'_2$	$-U_2$	X'_2	X_2	$-U'_1$	$-U_1$	X_1	A_0

If $\dim \mathfrak{e} > 1$ we obtain other subgroups E of $E'_{\emptyset}A \cong Sp(1)Sp(1)\mathbb{R}$ such that EN acts transitively on $A_{Sp(2, 1)}$. Such a group E is isomorphic to some group of the form $U(1)\mathbb{R}$, $U(1)U(1)\mathbb{R}$, $Sp(1)\mathbb{R}$, $Sp(1)U(1)\mathbb{R}$ or $Sp(1)Sp(1)\mathbb{R}$. However, the natural reductive decompositions defined by their actions do not provide new homogeneous Riemannian structures.

We have proved

Proposition 3.2. *The space $A_{Sp(2, 1)}$ admits the homogeneous quaternionic Kähler structures S (associated to each corresponding homogeneous description*

\hat{G}/H) given by

S	\hat{G}/H
0	$Sp(2,1)/(Sp(2) \times Sp(1))$
Table II $(\lambda, \mu \in \mathbb{R}^3)$	$E_{\lambda, \mu} N$ $(E_{\lambda, \mu} \cong \mathbb{R})$

where N is the nilpotent factor in the Iwasawa decomposition of $Sp(2,1)$. \square

3.3 The exceptional space $A_{G_{2(2)}} = G_{2(2)}/SO(4)$

3.3.1 The quaternion-Hermitian structure of $\mathfrak{a} + \mathfrak{n} \subset \mathfrak{g}_{2(2)}$

The non-compact real form $\mathfrak{g}_{2(2)}$ of the exceptional Lie algebra \mathfrak{g}_2 can be realized as a subalgebra of $\mathfrak{so}(4,3)$. Each element $X \in \mathfrak{g}_{2(2)}$ can be represented by a real matrix of the form

$$(3.6) \quad \begin{pmatrix} 0 & x_1 + u_1 & -x_2 - u_2 & 2u_3 & s_1 & y_1 + v_1 & y_2 - v_2 \\ -x_1 - u_1 & 0 & x_3 + u_3 & 2u_2 & y_1 - v_1 & s_2 & y_3 + v_3 \\ x_2 + u_2 & -x_3 - u_3 & 0 & 2u_1 & y_2 + v_2 & y_3 - v_3 & s_3 \\ -2u_3 & -2u_2 & -2u_1 & 0 & 2v_3 & 2v_2 & 2v_1 \\ s_1 & y_1 - v_1 & y_2 + v_2 & 2v_3 & 0 & x_1 - u_1 & -x_2 + u_2 \\ y_1 + v_1 & s_2 & y_3 - v_3 & 2v_2 & -x_1 + u_1 & 0 & x_3 - u_3 \\ y_2 - v_2 & y_3 + v_3 & s_3 & 2v_1 & x_2 - u_2 & -x_3 + u_3 & 0 \end{pmatrix},$$

with $s_1 + s_2 + s_3 = 0$. The involution τ of $\mathfrak{g}_{2(2)}$ given by $\tau(X) = -X^T$ defines the Cartan decomposition $\mathfrak{g}_{2(2)} = \mathfrak{k} + \mathfrak{p}$, where $\mathfrak{k} = (\mathfrak{so}(4) \oplus \mathfrak{so}(3)) \cap \mathfrak{g}_{2(2)} \cong \mathfrak{so}(4)$. The subspace \mathfrak{a} of \mathfrak{p} defined by

$$\mathfrak{a} = \left\{ \begin{pmatrix} 0 & 0 & 0 & 0 & s_1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & s_2 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & s_3 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ s_1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & s_2 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & s_3 & 0 & 0 & 0 & 0 \end{pmatrix} : s_1 + s_2 + s_3 = 0 \right\}$$

is a maximal \mathbb{R} -diagonalizable subalgebra of $\mathfrak{g}_{2(2)}$, and $Z_{\mathfrak{k}}(\mathfrak{a}) = \{0\}$. The elements A_1 and A_2 of \mathfrak{a} defined by $(s_1, s_2, s_3) = (1, -1, 0)$ and $(-2, 1, 1)$, respectively, generate \mathfrak{a} . Let f_j be the member of \mathfrak{a}^* whose value on the above matrix of \mathfrak{a} is s_j . Then $f_1 + f_2 + f_3 = 0$. We consider the (weak) order in \mathfrak{a}^* defined by choosing the element $A_0 \in \mathfrak{a}$ given by $(s_1, s_2, s_3) = (-1, -2, 3)$ and by saying that $f \in \mathfrak{a}^*$ is positive if $f(A_0) > 0$. Then, the sets of positive roots and simple roots of $(\mathfrak{g}_{2(2)}, \mathfrak{a})$ are $\Sigma^+ = \{f_1 - f_2, f_3 - f_1, f_3 - f_2, f_3, -f_2, -f_1\}$ and $\Pi = \{f_1 - f_2, -f_1\}$, respectively. Now, we will give the generators X_j, Y_j, U_j, V_j , $1 \leq j \leq 3$, of the root spaces \mathfrak{g}_f , $f \in \Sigma$. The matrix X in (3.6) represents

$$\begin{aligned} X_1 & \text{ if } x_1 = y_1 = \frac{1}{2}, & X_2 & \text{ if } x_2 = y_2 = \frac{1}{2}, & X_3 & \text{ if } x_3 = -y_3 = -\frac{1}{2}, \\ Y_1 & \text{ if } x_1 = -y_1 = -\frac{1}{2}, & Y_2 & \text{ if } x_2 = -y_2 = -\frac{1}{2}, & Y_3 & \text{ if } x_3 = y_3 = -\frac{1}{2}, \\ U_1 & \text{ if } u_1 = v_1 = -\frac{1}{2}, & U_2 & \text{ if } u_2 = -v_2 = \frac{1}{2}, & U_3 & \text{ if } u_3 = -v_3 = -\frac{1}{2}, \\ V_1 & \text{ if } u_1 = -v_1 = \frac{1}{2}, & V_2 & \text{ if } u_2 = v_2 = -\frac{1}{2}, & V_3 & \text{ if } u_3 = v_3 = \frac{1}{2}, \end{aligned}$$

with all the other values $x_j, y_j, u_j, v_j, s_j = 0$ for each of the 12 cases. We have $\mathfrak{g}_{f_1-f_2} = \langle X_1 \rangle$, $\mathfrak{g}_{-f_1+f_2} = \langle Y_1 \rangle$, $\mathfrak{g}_{f_3-f_1} = \langle X_2 \rangle$, $\mathfrak{g}_{-f_3+f_1} = \langle Y_2 \rangle$, $\mathfrak{g}_{f_3-f_2} = \langle X_3 \rangle$, $\mathfrak{g}_{-f_3+f_2} = \langle Y_3 \rangle$, $\mathfrak{g}_{f_3} = \langle U_1 \rangle$, $\mathfrak{g}_{-f_3} = \langle V_1 \rangle$, $\mathfrak{g}_{-f_2} = \langle U_2 \rangle$, $\mathfrak{g}_{f_2} = \langle V_2 \rangle$, $\mathfrak{g}_{-f_1} = \langle U_3 \rangle$, $\mathfrak{g}_{f_1} = \langle V_3 \rangle$. Then $\mathfrak{g}_{2(2)} = \mathfrak{a} + \sum_{f \in \Sigma} \mathfrak{g}_f$, and we have the Iwasawa decomposition $\mathfrak{g}_{2(2)} = \mathfrak{k} + \mathfrak{a} + \mathfrak{n}$, where $\mathfrak{n} = \sum_{f \in \Sigma^+} \mathfrak{g}_f = \langle X_1, X_2, X_3, U_1, U_2, U_3 \rangle$.

The elements E_1, E_2, E_3 of $\mathfrak{k} = \mathfrak{so}(4) \subset \mathfrak{g}_{2(2)}$, which are represented by the matrices

$$\begin{pmatrix} 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}, \begin{pmatrix} 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}, \begin{pmatrix} 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 \\ -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix},$$

respectively, satisfy $[E_1, E_2] = 2E_3$, $[E_2, E_3] = 2E_1$, $[E_3, E_1] = 2E_2$, and generate a compact ideal $\mathfrak{u} \cong \mathfrak{sp}(1)$ of \mathfrak{k} , such that the isotropy representation $\mathfrak{u} \rightarrow \mathfrak{gl}(\mathfrak{p})$ defines a quaternionic structure on $A_{G_{2(2)}}$. Through the isomorphisms $\mathfrak{p} \cong \mathfrak{g}_{2(2)} / \mathfrak{k} \cong \mathfrak{a} + \mathfrak{n}$, the action of each E_i on \mathfrak{p} defines the complex structure J_a ($a = 1, 2, 3$), acting on the elements A_1, A_2, X_j, U_j ($1 \leq j \leq 3$) of $\mathfrak{a} + \mathfrak{n}$ as follows.

	A_1	A_2	X_1	X_2	X_3	U_1	U_2	U_3
J_1	$-2X_1$	$3X_1+U_1$	$\frac{1}{2}A_1$	$\frac{1}{2}(X_3-U_3)$	$-\frac{1}{2}(X_2+U_2)$	$-\frac{3}{2}A_1-A_2$	$\frac{1}{2}(3X_3+U_3)$	$\frac{1}{2}(3X_2-U_2)$
J_2	X_2-U_2	$-3X_2+U_2$	$-\frac{1}{2}(X_3+U_3)$	$\frac{1}{2}(A_1+A_2)$	$\frac{1}{2}(X_1-U_1)$	$\frac{1}{2}(3X_3-U_3)$	$\frac{1}{2}(3A_1+A_2)$	$\frac{1}{2}(3X_1+U_1)$
J_3	$-X_3-U_3$	$2U_3$	$\frac{1}{2}(-X_2+U_2)$	$\frac{1}{2}(X_1+U_1)$	$A_1+\frac{1}{2}A_2$	$-\frac{1}{2}(3X_2+U_2)$	$\frac{1}{2}(-3X_1+U_1)$	$-\frac{1}{2}A_2$

We consider the scalar product $\langle \cdot, \cdot \rangle$ induced in $\mathfrak{a} + \mathfrak{n}$ by the isomorphisms $\mathfrak{p} \cong \mathfrak{a} + \mathfrak{n}$ and $\frac{1}{16}B|_{\mathfrak{p} \times \mathfrak{p}}$, where B is the Killing form of $\mathfrak{g}_{2(2)}$. Then, $\langle A_1, A_1 \rangle = 1$, $\langle A_2, A_2 \rangle = 3$, $\langle A_1, A_2 \rangle = -\frac{3}{2}$, $\langle X_j, X_j \rangle = \frac{1}{4}$, $\langle U_j, U_j \rangle = \frac{3}{4}$, and $(\mathfrak{a} + \mathfrak{n}, \langle \cdot, \cdot \rangle, J_1, J_2, J_3)$ is a quaternion-Hermitian vector space.

3.3.2 Homogeneous descriptions and structures of $A_{G_{2(2)}}$

The four standard parabolic subalgebras of $\mathfrak{g}_{2(2)}$ are parametrized by the subsets $\Pi, \emptyset, \Psi_1 = \{f_1 - f_2\}$ and $\Psi_2 = \{-f_1\}$ of Π . By using Theorems 2.2 and 2.3 we will obtain the homogeneous descriptions of $A_{G_{2(2)}}$.

The case $\Psi = \Pi$. We have $[\Psi] = \Sigma$, $\mathfrak{e}'_{\Pi} = \mathfrak{a}_{\Pi} = \mathfrak{n}_{\Pi} = \{0\}$, and the refined Langlands decomposition is $\mathfrak{p}_{\Pi} = \mathfrak{g}_{2(2)} + \{0\} + \{0\} + \{0\}$. Here one only gets the transitive action of the full isometry group $G_{2(2)}$ on $A_{G_{2(2)}}$, and we have the description of $A_{G_{2(2)}}$ as the symmetric space $G_{2(2)}/SO(4)$. The associated reductive decomposition is the Cartan decomposition $\mathfrak{g}_{2(2)} = \mathfrak{k} + \mathfrak{p}$, with $\mathfrak{k} \cong \mathfrak{so}(4)$, and the corresponding structure is $S = 0$.

The case $\Psi = \emptyset$. In this case, $\mathfrak{l}'_{\emptyset} = \{0\}$, $\mathfrak{e}'_{\emptyset} = Z_{\mathfrak{k}}(\mathfrak{a}) = \{0\}$, and $\mathfrak{a}_{\emptyset} = \mathfrak{a}$, and hence the refined Langlands decomposition of the minimal parabolic subalgebra is $\mathfrak{p}_{\emptyset} = \{0\} + \{0\} + \mathfrak{a} + \mathfrak{n}$. In this case, the only connected co-compact subgroup \hat{G} which acts transitively on $A_{G_{2(2)}}$ is provided by $E = A$, so

we have the description of $A_{G_{2(2)}}$ as the solvable Lie group $\hat{G} = AN$, where N is the nilpotent factor in the Iwasawa decomposition of $G_{2(2)}$. The associated reductive decomposition is $\mathfrak{a} + \mathfrak{n} = \{0\} + (\mathfrak{a} + \mathfrak{n})$ and, having in mind that in this case $\mathfrak{m} = \mathfrak{a} + \mathfrak{n}$, by using (2.4) we obtain that the corresponding homogeneous quaternionic Kähler structure S is given by Table III with $\delta = \varepsilon = 1$.

Table III	A_1	A_2	X_1	X_2	X_3	U_1	U_2	U_3
S_{A_1}	0	0	0	0	0	0	0	0
S_{A_2}	0	0	0	0	0	0	0	0
S_{X_1}	$-\delta 2X_1$	$3\delta X_1$	$\frac{1}{2}\delta A_1$	$\frac{1}{2}\delta X_3$	$-\frac{1}{2}\delta X_2$	0	$-\frac{1}{2}\delta U_3$	$\frac{1}{2}\delta U_2$
S_{X_2}	X_2	$-3X_2$	$-\frac{1}{2}X_3$	$\frac{1}{2}A_1 + \frac{1}{2}A_2$	$\frac{1}{2}X_1$	$\frac{1}{2}U_3$	0	$-\frac{1}{2}U_1$
S_{X_3}	$-X_3$	0	$-\frac{1}{2}X_2$	$\frac{1}{2}X_1$	$A_1 + \frac{1}{2}A_2$	$\frac{1}{2}U_2$	$-\frac{1}{2}U_1$	0
S_{U_1}	0	$-U_1$	0	$\frac{1}{2}U_3$	$\frac{1}{2}U_2$	$\frac{3}{2}A_1 + A_2$	$-\frac{3}{2}X_3 - U_3$	$-\frac{3}{2}X_2 + U_2$
S_{U_2}	$-U_2$	U_2	$-\frac{1}{2}U_3$	0	$-\frac{1}{2}U_1$	$\frac{3}{2}X_3 - U_3$	$\frac{3}{2}A_1 + \frac{1}{2}A_2$	$\frac{3}{2}X_1 + U_1$
S_{U_3}	εU_3	$-2\varepsilon U_3$	$-\frac{1}{2}\varepsilon U_2$	$-\frac{1}{2}\varepsilon U_1$	0	$\varepsilon(\frac{3}{2}X_2 + U_2)$	$\varepsilon(\frac{3}{2}X_1 - U_1)$	$\frac{1}{2}\varepsilon A_2$

The case $\Psi = \Psi_1$. Then $[\Psi_1] = \{\pm(f_1 - f_2)\}$, and $\mathfrak{p}_{\Psi_1} = \mathfrak{l}'_{\Psi_1} + \{0\} + \mathfrak{a}_{\Psi_1} + \mathfrak{n}_{\Psi_1}$, where $\mathfrak{a}_{\Psi_1} = \langle 3A_1 + 2A_2 \rangle$, $\mathfrak{n}_{\Psi_1} = \sum_{f \in \Sigma^+, f \neq f_1 - f_2} \mathfrak{g}_f = \langle X_2, X_3, U_1, U_2, U_3 \rangle$, and

$$\mathfrak{l}'_{\Psi_1} = \left\{ \begin{pmatrix} 0 & x & 0 & 0 & s & y & 0 & 0 \\ -x & 0 & 0 & 0 & 0 & y & -s & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ s & y & 0 & 0 & 0 & x & 0 & 0 \\ y & -s & 0 & 0 & -x & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix} \right\} = \langle A_1, X_1, Y_1 \rangle \cong \mathfrak{sl}(2, \mathbb{R}).$$

The connected Lie subgroup $E = A_{\Psi_1} \cong \mathbb{R}$ of P_{Ψ_1} with Lie algebra \mathfrak{a}_{Ψ_1} is the only possible choice to get a co-compact subgroup $\hat{G} = LEN_{\Psi_1}$ of $G_{2(2)}$, with $L = L'_{\Psi_1} \cong Sl(2, \mathbb{R})$, and \hat{G} acts transitively on $A_{G_{2(2)}}$. The isotropy algebra is $\mathfrak{h} = \hat{\mathfrak{g}} \cap \mathfrak{k} = \mathfrak{l}'_{\Psi_1} \cap \mathfrak{k} = \langle (X_1)_{\mathfrak{k}} \rangle \cong \mathfrak{so}(2)$. We have the reductive decomposition $\hat{\mathfrak{g}} = \mathfrak{h} + \mathfrak{m}$, where $\hat{\mathfrak{g}} = \mathfrak{p}_{\Psi_1}$, and $\mathfrak{m} = \langle A_1, A_2, (X_1)_{\mathfrak{p}}, X_2, X_3, U_1, U_2, U_3 \rangle$, which is associated to the description $A_{G_{2(2)}} \equiv Sl(2, \mathbb{R})\mathbb{R}N_{\Psi_1}/SO(2)$. By using (2.4), we obtain that the corresponding homogeneous quaternionic Kähler structure S is given at o by Table III with $\delta = 0$, $\varepsilon = 1$.

The case $\Psi = \Psi_2$. Then $[\Psi_2] = \{\pm f_1\}$ and $\mathfrak{p}_{\Psi_2} = \mathfrak{l}'_{\Psi_2} + \{0\} + \mathfrak{a}_{\Psi_2} + \mathfrak{n}_{\Psi_2}$, where $\mathfrak{a}_{\Psi_2} = \langle 2A_1 + A_2 \rangle$, $\mathfrak{n}_{\Psi_2} = \sum_{f \in \Sigma^+, f \neq -f_1} \mathfrak{g}_f = \langle X_1, X_2, X_3, U_1, U_2 \rangle$, and

$$\mathfrak{l}'_{\Psi_2} = \left\{ \begin{pmatrix} 0 & 0 & 0 & 2u & -2s & 0 & 0 & 0 \\ 0 & 0 & u & 0 & 0 & s & v & 0 \\ 0 & -u & 0 & 0 & 0 & -v & s & 0 \\ -2u & 0 & 0 & 0 & 2v & 0 & 0 & 0 \\ -2s & 0 & 0 & 2v & 0 & 0 & 0 & 0 \\ 0 & s & -v & 0 & 0 & 0 & -u & 0 \\ 0 & v & s & 0 & 0 & u & 0 & 0 \end{pmatrix} \right\} = \langle A_2, U_3, V_3 \rangle \cong \mathfrak{sl}(2, \mathbb{R}).$$

In this case we also have only one possible choice for E as the connected subgroup $A_{\Psi_2} \cong \mathbb{R}$ of P_{Ψ_2} whose Lie algebra is \mathfrak{a}_{Ψ_2} . So, we get a co-compact subgroup $\hat{G} = LEN_{\Psi_2}$ of $G_{2(2)}$, with $L = L'_{\Psi_2} \cong Sl(2, \mathbb{R})$, which acts transitively on $A_{G_{2(2)}}$. This gives $A_{G_{2(2)}} \equiv Sl(2, \mathbb{R})\mathbb{R}N_{\Psi_2}/SO(2)$, whose natural associated

reductive decomposition is $\hat{\mathfrak{g}} = \mathfrak{h} + \mathfrak{m}$, where $\hat{\mathfrak{g}} = \mathfrak{p}_{\Psi_2}$, $\mathfrak{h} = \hat{\mathfrak{g}} \cap \mathfrak{k} = \mathfrak{l}'_{\Psi_2} \cap \mathfrak{k} = \langle (U_3)_{\mathfrak{k}} \rangle \cong \mathfrak{so}(2)$ and $\mathfrak{m} = \langle A_1, A_2, X_1, X_2, X_3, U_1, U_2, (U_3)_{\mathfrak{p}} \rangle$. By (2.4), we have that the corresponding structure S is given at o by Table III with $\delta = 1$, $\varepsilon = 0$.

We have

Proposition 3.3. *The space $A_{G_{2(2)}}$ admits the homogeneous quaternionic Kähler structures S (associated to each corresponding homogeneous description \hat{G}/H) given by*

S	\hat{G}/H
0	$G_{2(2)}/SO(4)$
Table III, $\delta = \varepsilon = 1$	AN ($A \cong \mathbb{R}^2$)
Table III, $\delta = 0$, $\varepsilon = 1$	$Sl(2, \mathbb{R})\mathbb{R}N_{\Psi_1}/SO(2)$
Table III, $\delta = 1$, $\varepsilon = 0$	$Sl(2, \mathbb{R})\mathbb{R}N_{\Psi_2}/SO(2)$

where N is the nilpotent factor in the Iwasawa decomposition $G_{2(2)} = SO(4)AN$ and Ψ_1 and Ψ_2 are the subsets of simple restricted roots defining the intermediate parabolic subalgebras of $(\mathfrak{g}_{2(2)}, \mathfrak{a})$. \square

3.4 Types of homogeneous quaternionic Kähler structures

For each parabolic subgroup P_{Ψ} of the full connected isometry group G of each non-compact quaternion-Kähler symmetric space M of dimension 8, the subgroups \hat{G} of G acting transitively on M are of the form $\hat{G} = LEN_{\Psi}$ as in Theorem 2.3, where $L = L'_{\Psi}$. Moreover, E is a connected closed subgroup of $E'_{\Psi}A_{\Psi}$ such that the projection of its Lie algebra $\mathfrak{e} \subset \mathfrak{e}'_{\Psi} + \mathfrak{a}_{\Psi}$ to \mathfrak{a}_{Ψ} is surjective. If this projection is an isomorphism we say that E is *minimal*. Then we have

Theorem 3.4. *Let $G = KAN$ be the Iwasawa decomposition of each of the groups $SU(2, 2)$, $Sp(2, 1)$, and $G_{2(2)}$. The homogeneous descriptions of each eight-dimensional non-compact quaternion-Kähler symmetric space are of the form \hat{G}/H , where $\hat{G} = L'_{\Psi}EN_{\Psi}$, with L'_{Ψ} either non-compact simple or trivial and N_{Ψ} nilpotent. If E is minimal, then it is simply-connected and abelian, and in this case the corresponding types of homogeneous quaternionic Kähler structures are given in the following table (where the figure on the fifth column, if any, stands for the number of parameters of the corresponding n -parametric family of homogeneous quaternionic Kähler structures).*

IV	G/K	Ψ	$L'_{\Psi}EN_{\Psi}/H$	$\dim E$	n	type
	$A_{SU(2,2)}$	Π	$SU(2, 2)/S(U(2) \times U(2))$	0		$\{0\}$
		\emptyset	$E_{\lambda, \mu}N$ ($\lambda, \mu \in \mathbb{R}$)	2	2	\mathcal{QK}_{12345}
		Ψ_1	$Sl(2, \mathbb{C})A_{\Psi_1}N_{\Psi_1}/SU(2)$	1		\mathcal{QK}_{135}
		Ψ_2	$SU(1, 1)E_{\lambda}N_{\Psi_2}/U(1)$	1	1	\mathcal{QK}_{12345}
	$A_{Sp(2,1)}$	Π	$Sp(2, 1)/(Sp(2) \times Sp(1))$	0		$\{0\}$
		\emptyset	$E_{\lambda, \mu}N$ ($\lambda, \mu \in \mathbb{R}^3 \setminus \{0\}$)	1	6	\mathcal{QK}_{12345}
		\emptyset	$E_{0, \mu}N$ ($\mu \in \mathbb{R}^3 \setminus \{0\}$)	1	3	\mathcal{QK}_{1345}
		\emptyset	$AN = E_{0,0}N$	1		\mathcal{QK}_{134}
	$A_{G_{2(2)}}$	Π	$G_{2(2)}/SO(4)$	0		$\{0\}$
		\emptyset	AN	2		\mathcal{QK}_{12345}
		Ψ_1	$Sl(2, \mathbb{R})A_{\Psi_1}N_{\Psi_1}/SO(2)$	1		\mathcal{QK}_{12345}
		Ψ_2	$Sl(2, \mathbb{R})A_{\Psi_2}N_{\Psi_2}/SO(2)$	1		\mathcal{QK}_{12345}

Proof. In order to determine the types of non-zero homogeneous quaternionic Kähler structures (see Theorem 1.3) on the three spaces, we obtain for each structure the forms θ^i satisfying equations (1.3), by using the formula (2.6).

Case 1. The space $A_{SU(2,2)}$. The values of the forms θ^i at $o \in A_{SU(2,2)}$ corresponding to the homogeneous quaternionic Kähler structure S given in Table I are as follows

	A_1	A_2	U_1	U_2	X_1	X'_1	X_2	X'_2
θ^1	2λ	2μ	1	$-\delta$	0	0	0	0
θ^2	0	0	0	0	-2	0	2ε	0
θ^3	0	0	0	0	0	-2	0	2ε

If $\Psi = \emptyset$, ($\delta = \varepsilon = 1$), the structure $S = S^{\lambda, \mu}$ ($\lambda, \mu \in \mathbb{R}$) decomposes as $S = \Theta + T$, with $\Theta \in \mathcal{QK}_{12}$ (i.e., $\Theta_X Y = \frac{1}{2} \sum_{a=1}^3 \theta^a(X) J_a Y$), and $T \in \mathcal{QK}_{345}$. The condition for the tensor Θ to be in \mathcal{QK}_1 is $\theta^a = \theta \circ J_a$, $a = 1, 2, 3$, for some 1-form θ ; and this condition is not satisfied. In turn, $\Theta \in \mathcal{QK}_2$ if $\sum_{a=1}^3 \theta^a \circ J_a = 0$, which is not satisfied for instance by A_1 . Hence $\Theta \in \mathcal{QK}_{12} \setminus (\mathcal{QK}_1 \cup \mathcal{QK}_2)$. The relevant trace of T is $c_{12}(T) = \langle \frac{7}{2} A_1 + \frac{3}{2} A_2 + \lambda U_1 - \mu U_2, \cdot \rangle$. On the other hand, since $\dim M = 8$, the form defining the \mathcal{QK}_3 -component of T is given by $\vartheta = \frac{1}{10} c_{12}$. Hence $\vartheta = \frac{1}{20} \langle 7A_1 + 3A_2 + 2\lambda U_1 - 2\mu U_2, \cdot \rangle$, so the \mathcal{QK}_3 -component of T does not vanish either. Consider now the operator $F: \hat{\mathcal{V}} \rightarrow \hat{\mathcal{V}}$ defined by $F(T)_{XYZ} = T_{YZX} - T_{ZYX} + \sum_{a=1}^3 (T_{J_a Y J_a Z X} - T_{J_a Z J_a Y X})$, with eigenvalues 2 and -4 and respective eigenspaces \mathcal{QK}_{34} and \mathcal{QK}_5 (see Theorem 1.3). Take the tensor in \mathcal{QK}_3 given by $T_{XYZ}^\vartheta = \langle X, Y \rangle \vartheta(Z) - \langle X, Z \rangle \vartheta(Y) + \sum_{a=1}^3 (\langle X, J_a Y \rangle \vartheta(J_a Z) - \langle X, J_a Z \rangle \vartheta(J_a Y))$. Then $T - T^\vartheta \in \mathcal{QK}_{45}$ and we get $F(T - T^\vartheta)_{XYZ} = F(T)_{XYZ} - 2T_{XYZ}^\vartheta$. In particular, we obtain $(T - T^\vartheta)_{U_2 X'_1 X_2} = -\frac{17}{10}$ and $F(T - T^\vartheta)_{U_2 X'_1 X_2} = \frac{18}{5}$, so $T - T^\vartheta \in \mathcal{QK}_{45} \setminus (\mathcal{QK}_4 \cup \mathcal{QK}_5)$. Hence the tensor S has a non-zero component in each subspace \mathcal{QK}_j , $j = 1, \dots, 5$.

If $\Psi = \Psi_1$, ($\lambda = \mu = \varepsilon = 0$, $\delta = 1$), we have $S = \Theta + T$, where $\Theta \in \mathcal{QK}_1$, with corresponding 1-form $\theta = \langle A_1 + A_2, \cdot \rangle$, and T having non-zero \mathcal{QK}_3 -component, with corresponding 1-form $\vartheta = \frac{1}{4} \langle A_1 + A_2, \cdot \rangle$. Moreover, we get $(T - T^\vartheta)_{U_1 X'_2 X_1} = -\frac{3}{2}$ and $F(T - T^\vartheta)_{U_1 X'_2 X_1} = 6 = -4(T - T^\vartheta)_{U_1 X'_2 X_1}$. So $T - T^\vartheta \notin \mathcal{QK}_4$. As a computation with Maple shows, this happens for any choice of vectors. Hence $S \in \mathcal{QK}_{135}$.

If $\Psi = \Psi_2$, ($\mu = \delta = 0$, $\varepsilon = 1$), then $S = S^\lambda$ ($\lambda \in \mathbb{R}$) decomposes as $S = \Theta + T$ with $\Theta \in \mathcal{QK}_{12} \setminus (\mathcal{QK}_1 \cup \mathcal{QK}_2)$. Since moreover $\vartheta = \frac{1}{20} \langle 7A_1 + 2\lambda U_1, \cdot \rangle$, $(T - T^\vartheta)_{X'_1 U_2 X_2} = -\frac{13}{20}$, and $F(T - T^\vartheta)_{X'_1 U_2 X_2} = \frac{6}{5}$, the tensor S has a non-zero component in each subspace \mathcal{QK}_j , $j = 1, \dots, 5$.

Case 2. The space $A_{Sp(2,1)}$. The forms θ^i corresponding to the homogeneous quaternionic Kähler structure $S = S^{\lambda, \mu}$ ($\lambda, \mu \in \mathbb{R}^3$) given in Table II have the following values at $o \in A_{Sp(2,1)}$

	A_0	X_1	U_1	U'_1	X_2	X'_2	U_2	U'_2
θ^1	$2\lambda_1$	2	0	0	0	0	0	0
θ^2	$-2\lambda_2$	0	2	0	0	0	0	0
θ^3	$-2\lambda_3$	0	0	2	0	0	0	0

Equations (1.4) are satisfied, with the following nonzero values of θ^i at o : $\theta^1(A_0) = 2\lambda_1$, $\theta^2(A_0) = -2\lambda_2$, $\theta^3(A_0) = -2\lambda_3$, $\theta^1(X_1) = \theta^2(U_1) = \theta^3(U'_1) = 2$. We have $S = \Theta + T$, with $\Theta \in \mathcal{QK}_{12} \setminus \mathcal{QK}_1 \cup \mathcal{QK}_2$, except for $\lambda_1 = \lambda_2 = \lambda_3 = 0$. In this case $\Theta \in \mathcal{QK}_1$, with corresponding 1-form $\theta = 2\langle A_0, \cdot \rangle$. As for the specific type of T inside \mathcal{QK}_{345} , we first have that $\vartheta = \frac{1}{10}\langle 7A_0 + \lambda_1 X_1 - \lambda_2 U_1 - \lambda_3 U'_1, \cdot \rangle$, so the \mathcal{QK}_3 -component of T does not vanish. To find the \mathcal{QK}_{45} -component, we have for example $(T - T^\vartheta)_{U'_2 X_2 U'_1} = -\frac{3}{10}$ and $F(T - T^\vartheta)_{U'_2 X_2 U'_1} = -\frac{3}{5}$, so $T - T^\vartheta$ is never in \mathcal{QK}_5 . Suppose now that at least one of the parameters $\lambda_i, \mu_i, i = 1, 2, 3$, is non-zero. Computing now as before, we get for instance for $\mu_2 \neq 0$ that $(T - T^\vartheta)_{A_0 X_2 U_2} = -\mu_2$ and $F(T - T^\vartheta)_{A_0 X_2 U_2} = 2\mu_2$, which is not equal to $2(T - T^\vartheta)_{A_0 X_2 U_2}$. This also happens for the other parameters. Hence $T \in \mathcal{QK}_{345}$. If finally $\lambda_i, \mu_i, i = 1, 2, 3$, are all zero, a computation with Maple shows that $T \in \mathcal{QK}_{34}$ (see also [6, Prop. 5.3]).

Case 3. The space $A_{G_{2(2)}}$. The values of the forms θ^i at $o \in A_{G_{2(2)}}$ corresponding to the homogeneous quaternionic Kähler structure S given in Table III are the following

	A_1	A_2	X_1	X_2	X_3	U_1	U_2	U_3
θ^1	0	0	$\frac{\delta}{2}$	0	0	$-\frac{3}{2}$	0	0
θ^2	0	0	0	$\frac{1}{2}$	0	0	$\frac{3}{2}$	0
θ^3	0	0	0	0	$\frac{1}{2}$	0	0	$-\frac{3}{2}$

In any of the three cases, that is, $\Psi = \emptyset$, ($\delta = \varepsilon = 1$), $\Psi = \Psi_1$, ($\delta = 0, \varepsilon = 1$), and $\Psi = \Psi_2$, ($\delta = 1, \varepsilon = 0$), we have $S = \Theta + T$, with $\Theta \in \mathcal{QK}_{12} \setminus (\mathcal{QK}_1 \cup \mathcal{QK}_2)$. Since moreover one has

Ψ	ϑ	$(T - T^\vartheta)_{XYZ}$	$F(T - T^\vartheta)_{XYZ}$	vectors XYZ
\emptyset	$\frac{1}{30}\langle 21A_1 + 11A_2, \cdot \rangle$	$\frac{1}{20}$	$-\frac{29}{160}$	$X_1 X_2 X_3$
Ψ_1	$\frac{11}{60}\langle 3A_1 + 2A_2, \cdot \rangle$	$\frac{49}{160}$	$-\frac{29}{80}$	$X_2 X_3 X_1$
Ψ_2	$\frac{7}{20}\langle 2A_1 + A_2, \cdot \rangle$	$\frac{1}{20}$	$-\frac{73}{320}$	$X_1 X_2 X_3$

it follows that the tensor S has a non-zero component in each primitive subspace \mathcal{QK}_j , $j = 1, \dots, 5$. \square

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